

Hadron properties in dense nuclear matter

Gy. Wolf,

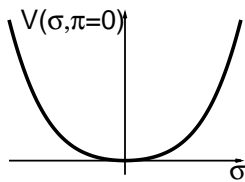
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In collaboration with M. Lutz, B. Friman and T. Hatsuda

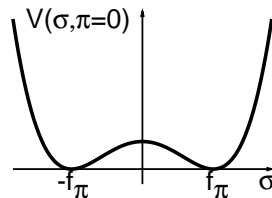
- Chiral symmetry restoration and the ρ and A_1 mesons
- Spectral function of vector mesons in matter
- QMD
- $A_1 - \rho$ mixing
- Summary

Chiral Symmetry

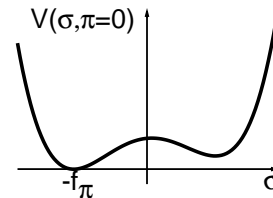
- for massless fermions: helicity
spin has the same or opposite direction to the momentum, right-handed or left-handed particles
- $m_u \approx 4\text{MeV}$, $m_d \approx 7\text{MeV}$, $m_s \approx 150\text{MeV}$
 $m_q \ll m_p$, QCD is approximately chiral symmetric
known before QCD (phenomenology of the weak interaction)
- The vacuum symmetric \Rightarrow parity doublets
not symmetric \Rightarrow SSB $\Rightarrow \langle \bar{q}q \rangle$ condensates in the vacuum



symmetry



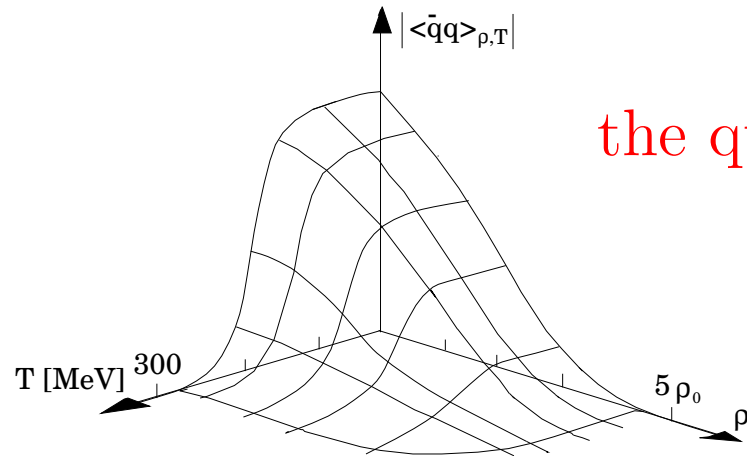
spontaneous



explicit breaking

Quark condensate

order parameter of the chiral phase transition: $\langle \bar{q}q \rangle$

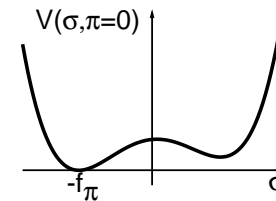


the quark condensate as a function of T and ρ

$$\frac{\langle \bar{q}q \rangle_{\rho_0}}{\langle \bar{q}q \rangle_{vac}} \approx 0.7$$

Proposed signals: disoriented chiral condensate (DCC) and change of hadron masses

- DCC: Large fluctuations in pion charges
coherent decay of the false vacuum

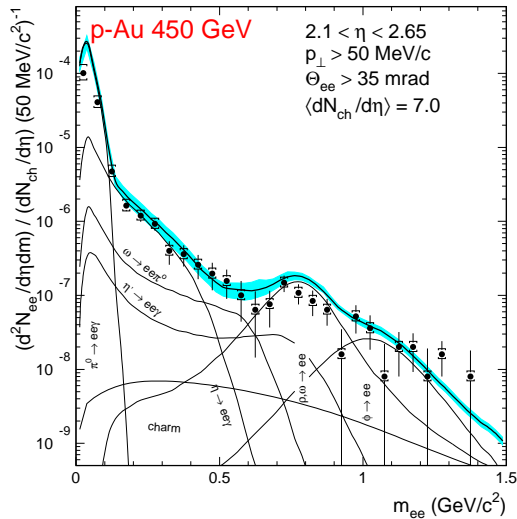
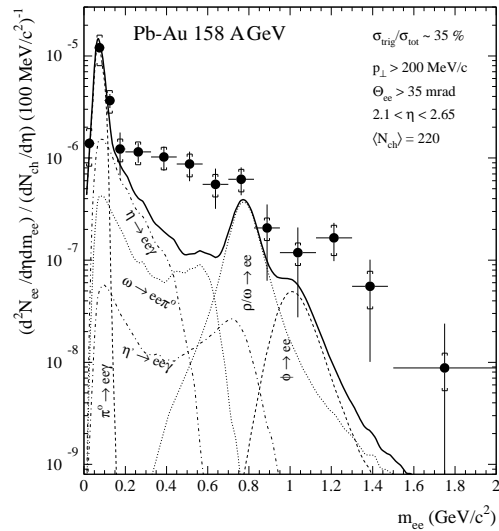
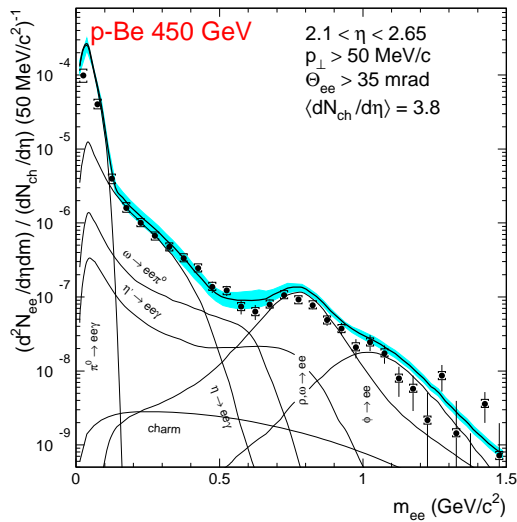


- Hadron masses: light vector meson masses shifted down
 \Rightarrow dilepton production

Chiral symmetry and vector mesons

- vector mesons can be observed by their dileptonic decay
- Chiral symmetry requires: in restored phase the spectral functions of parity partners (ρ and A_1) are the same.
- Two simplified scenarios for chiral restoration
 - masses become equal
 - spectral functions mix
- What do we know about the masses
 - A_1 : nothing, but in the vacuum
 - ρ : QCD sum rules (mass shifts)
 - hadronic models
- Study spectral functions

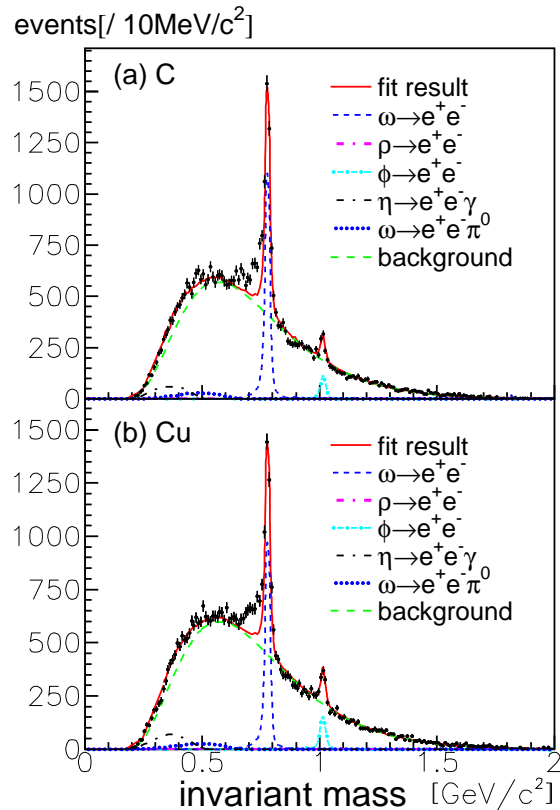
CERES data



G. Agakichiev *et al.*
 Eur. Phys. J. C4 (1998) 231

G. Agakichiev *et al.*
 Phys. Lett. B422 (1998) 405

p + A 12 GeV at KEK

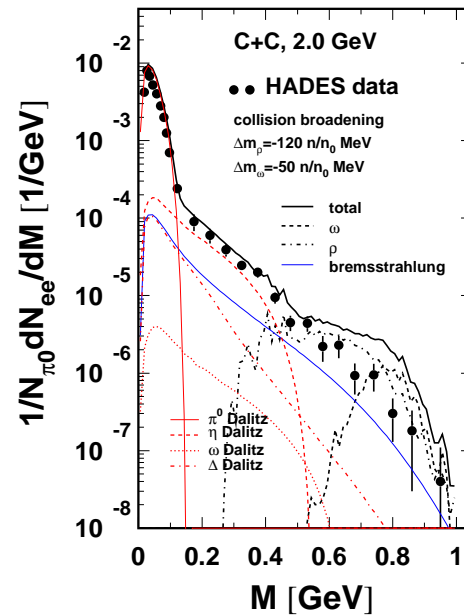
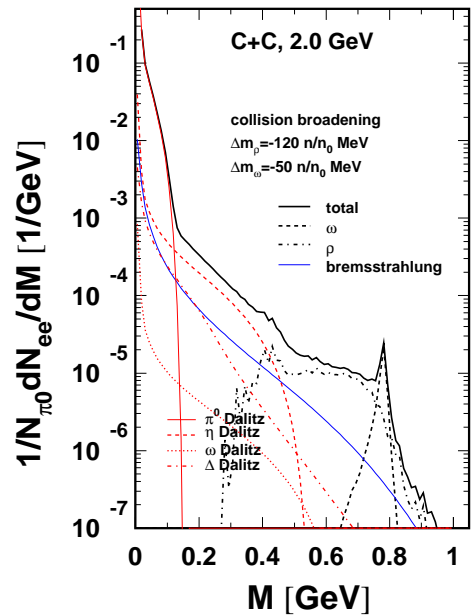
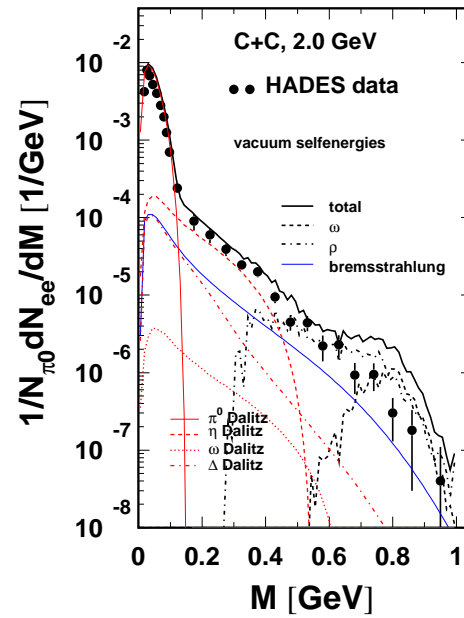
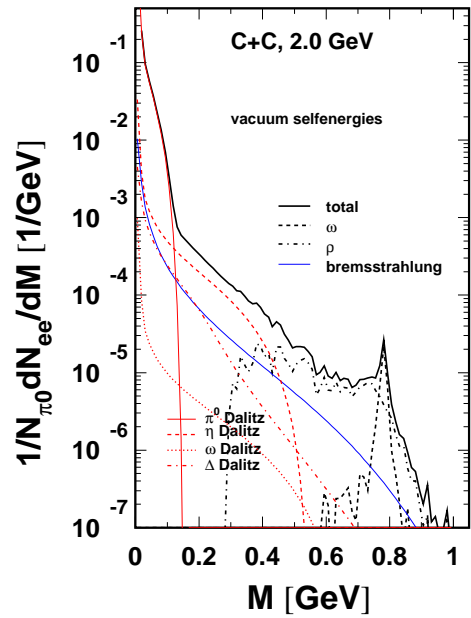


Naruki et al., Phys.Rev.Lett. 96 (2006) 092301

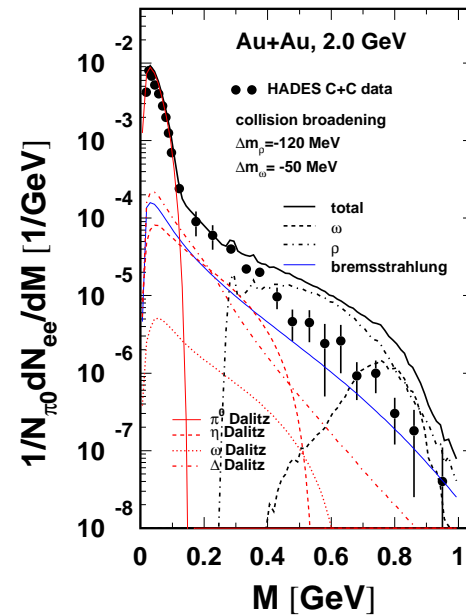
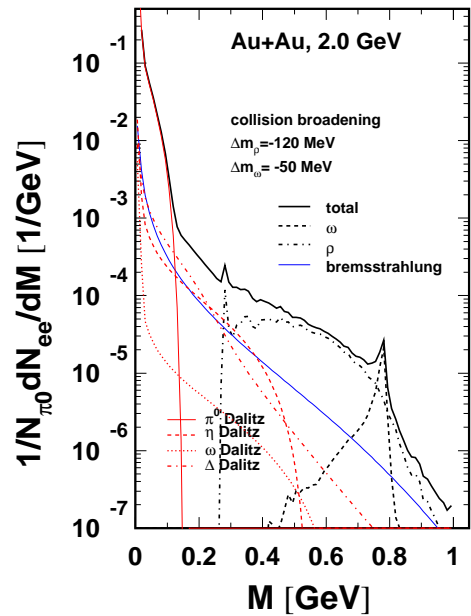
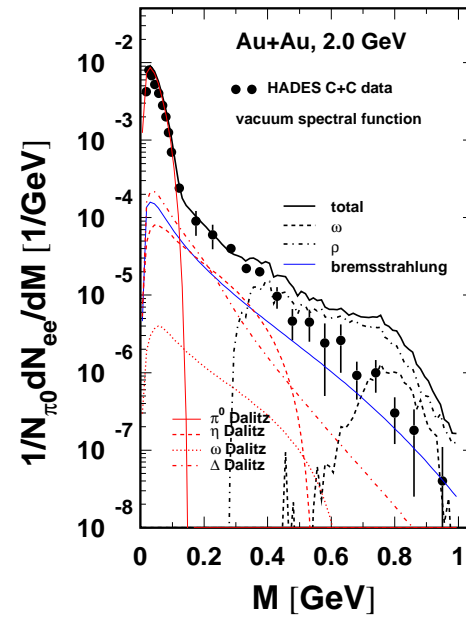
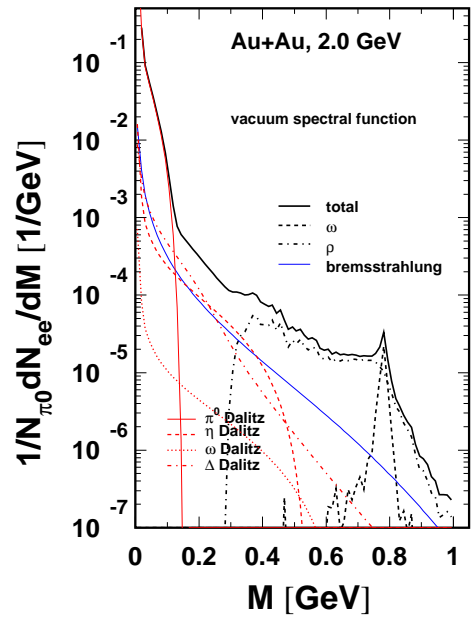
excess below the ω -mass, cannot be described by free cocktail

$$m_v = m_{v_0} * (1 - 0.09\rho/\rho_0), \quad v : \rho, \omega$$

C + C 2 GeV



Au + Au 2 GeV



Self energy of vector mesons

- low density expansion

$$\langle n.m. | \mathcal{O} | n.m. \rangle = \langle 0 | \mathcal{O} | 0 \rangle + \rho \langle N | \mathcal{O} | N \rangle$$

- Modification of vector mesons in nuclear matter (after LSZ reduction)

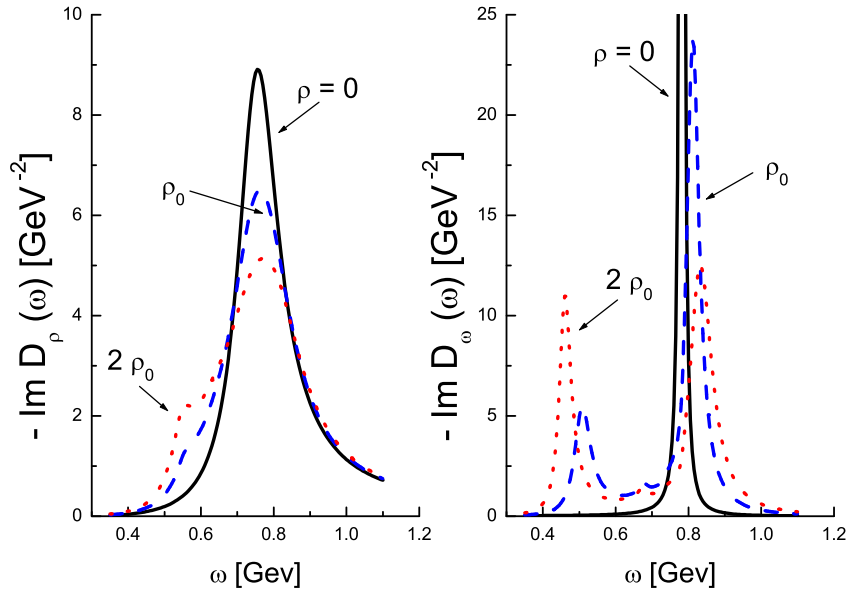
$$\Delta\Pi = -4\pi\rho\left(1 + \frac{m_V}{m_N}\right)\langle NV | NV \rangle$$

$\langle NV | NV \rangle$: vector mesons' forward scattering amplitude off the nucleon

- goal: $T_{\omega N \rightarrow \omega N}$ and $T_{\rho N \rightarrow \rho N}$
 - not measurable
 - measurable: $T_{\pi N \rightarrow V N}$ and $T_{\pi N \rightarrow \pi N}$

\Rightarrow coupled channels

Vector mesons in nuclear matter

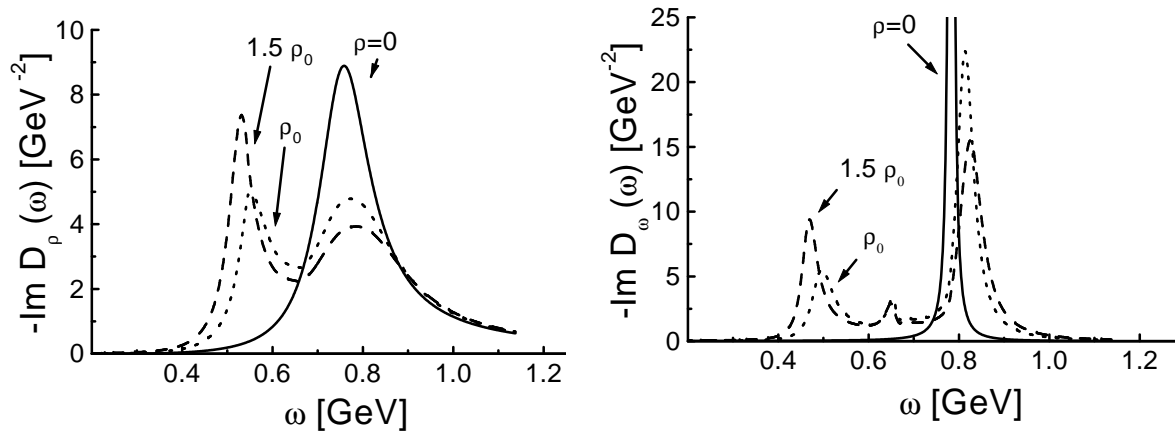


no mass shift, no simple broadening
but extra peaks: resonance-hole states
 ρ spectrum is model-dependent (ω stable)

no clean signal for chiral restoration

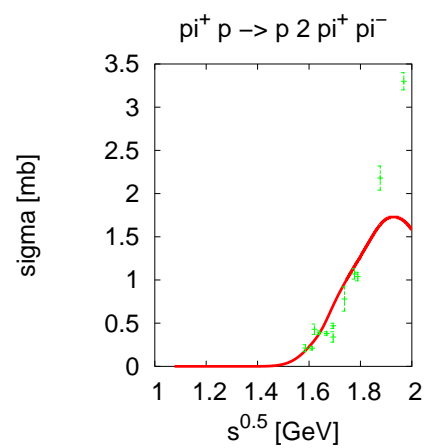
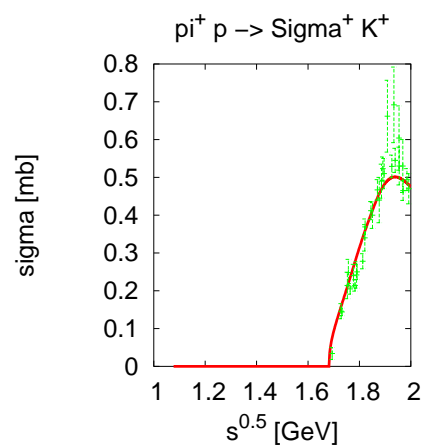
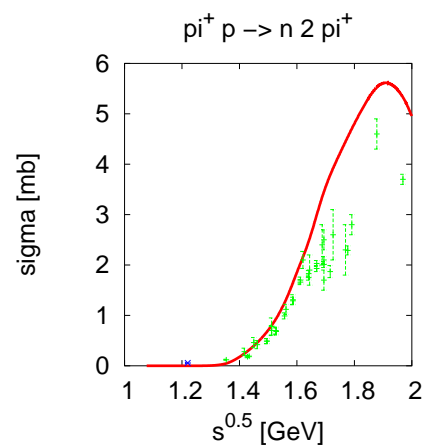
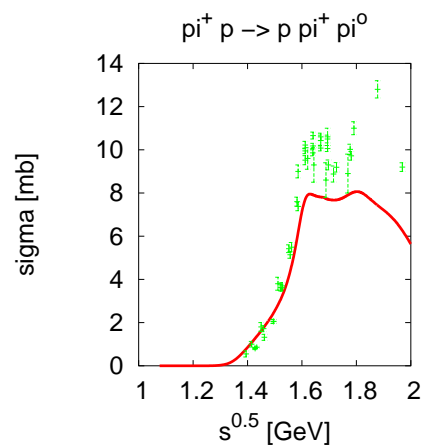
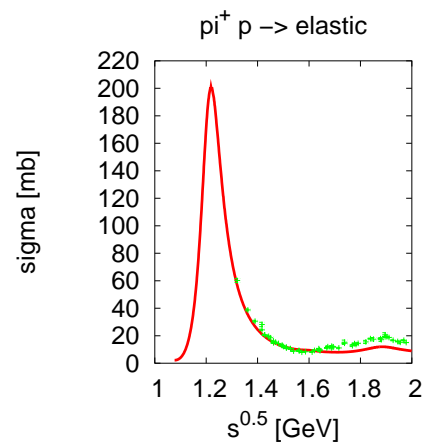
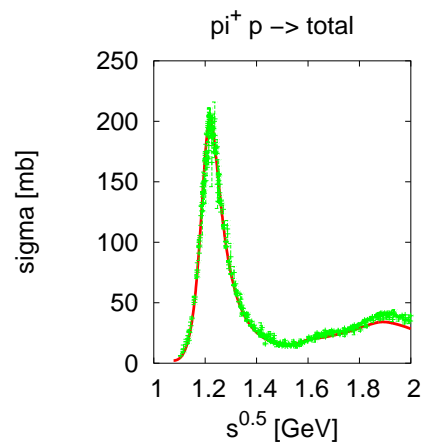
M. Lutz, Gy. Wolf, B. Friman; Nucl Phys. A706 (2002) 431.

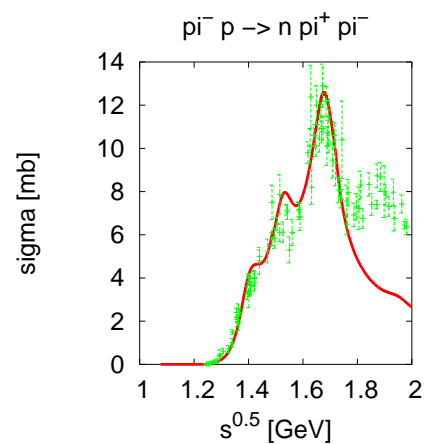
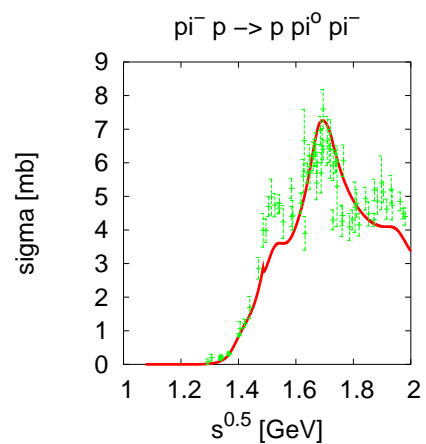
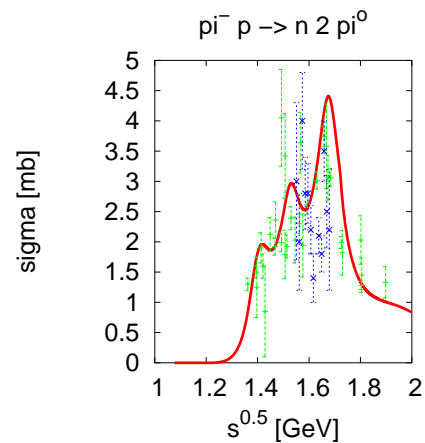
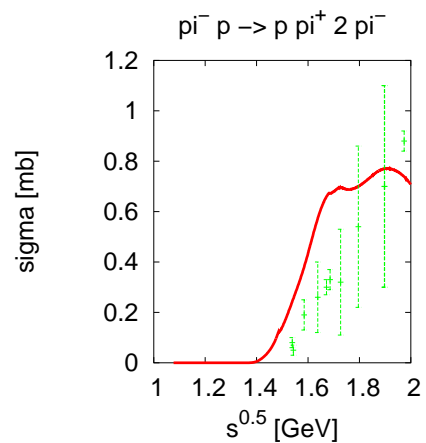
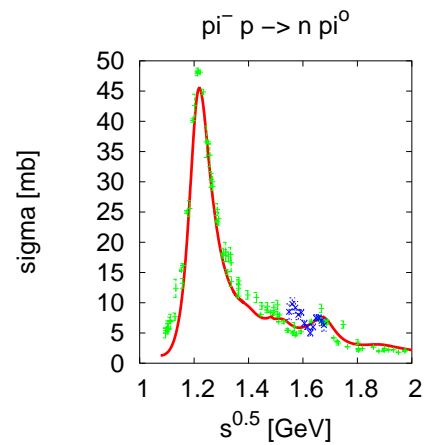
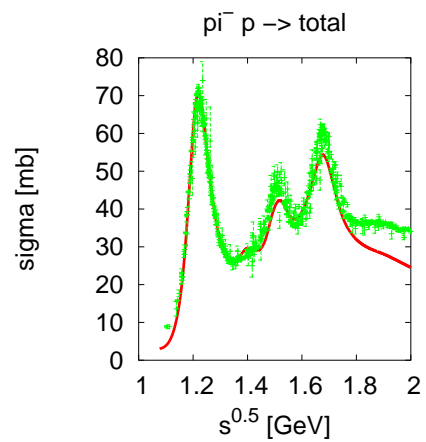
our old results:

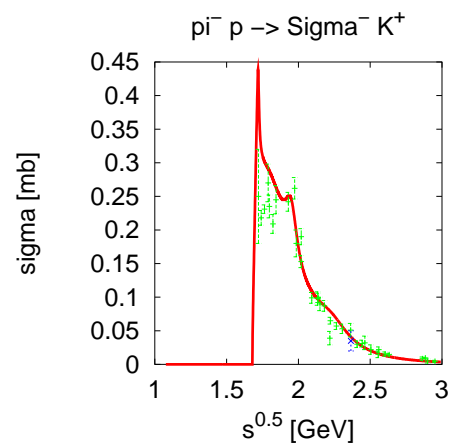
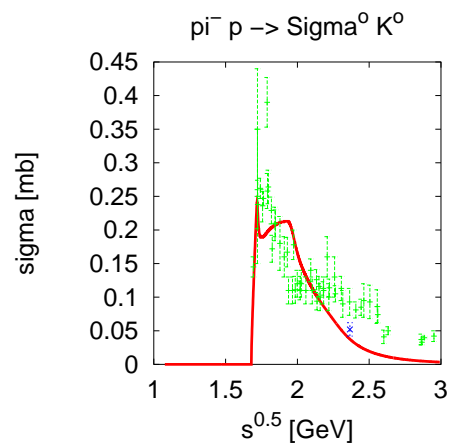
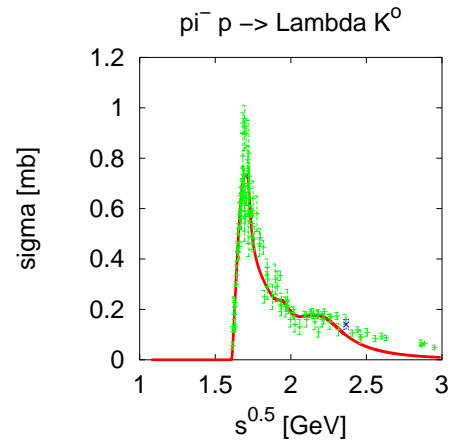
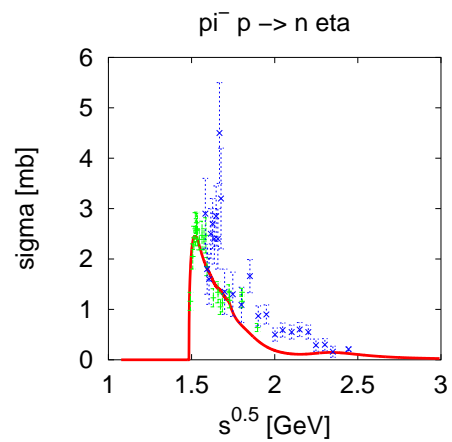
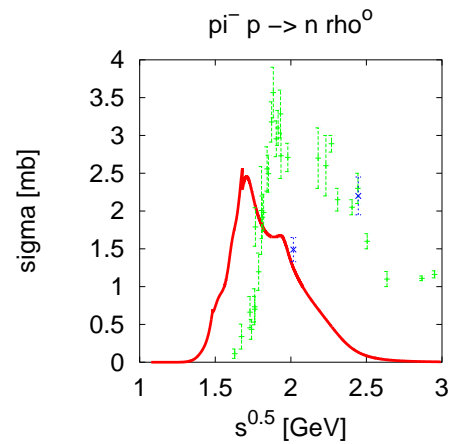
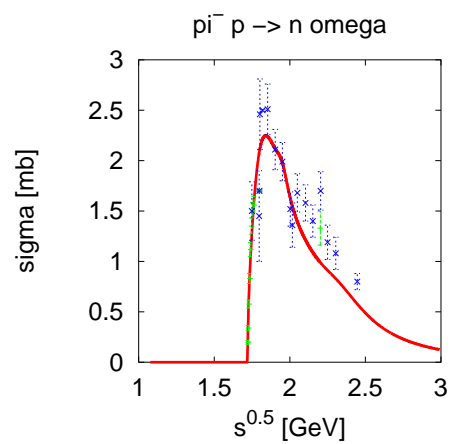


QMD

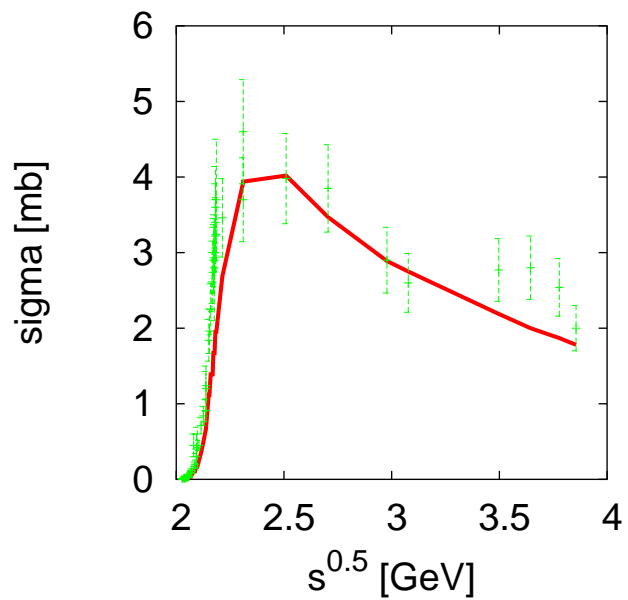
- Momentum dependent force between the particles
Model describes rather well the flow data at 400 MeV Au+Au
- 24 baryon resonances + Λ and Σ baryons
 $\pi, \eta, \sigma, \rho, \omega$ and kaons
- Inelastic reactions:
 $NN \leftrightarrow NR, \quad MN \leftrightarrow R, \quad \pi\pi \leftrightarrow \rho, \quad \omega \leftrightarrow \rho\pi$
- Cross sections are resonance dominated, and added incoherently
- resonance parameters are fitted to π -nucleon scattering
- resonance production cross sections are fitted to nucleon-nucleon cross sections



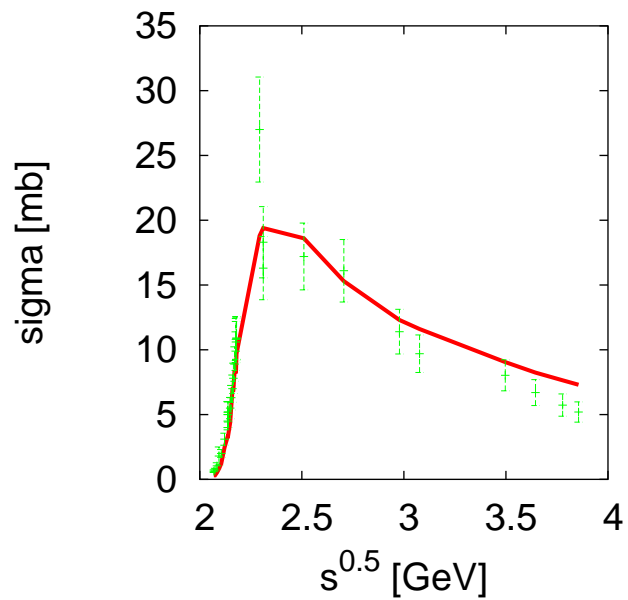




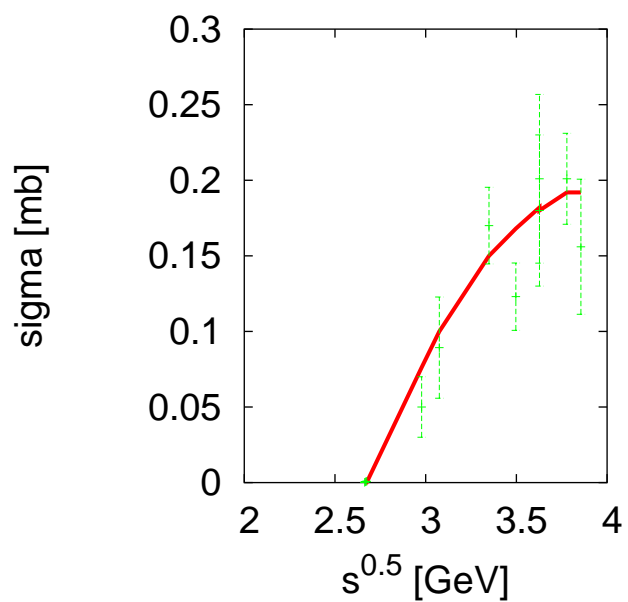
pp→pp pi⁰



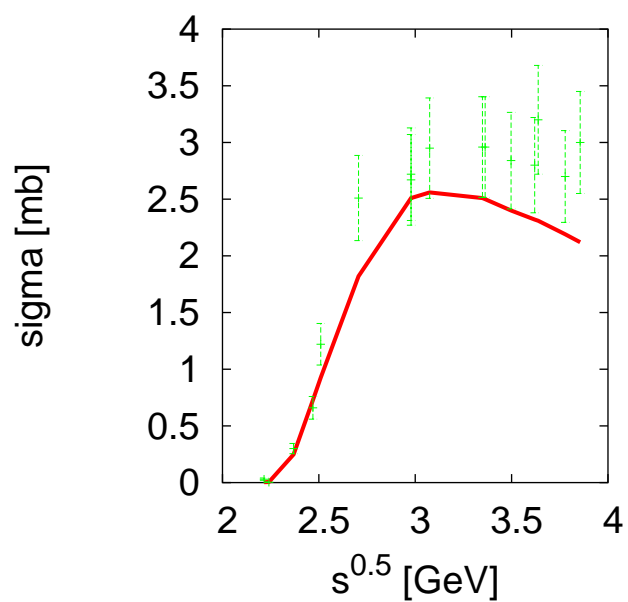
pp→pn pi⁺

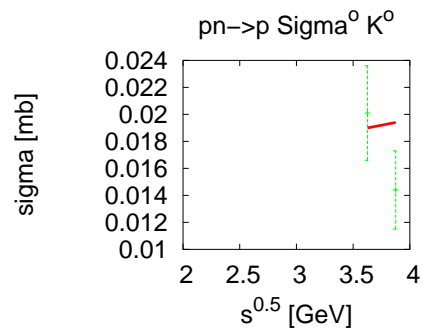
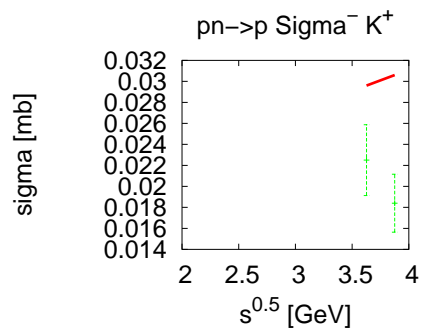
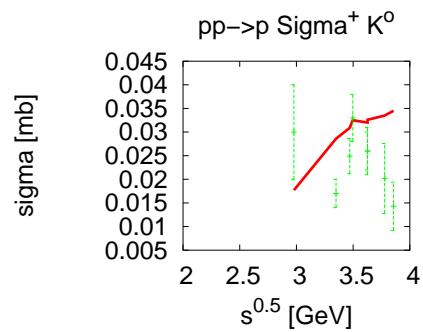
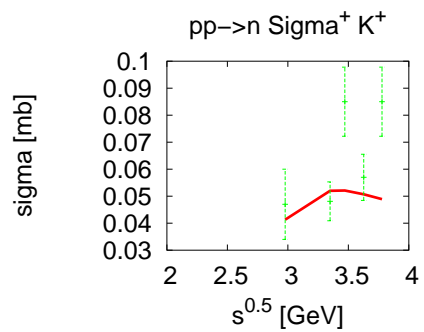
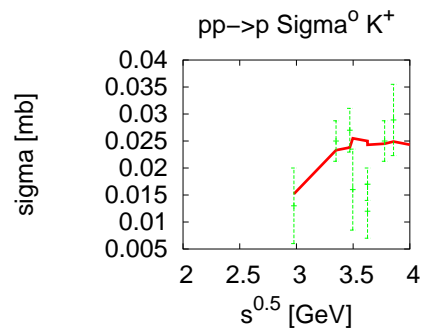
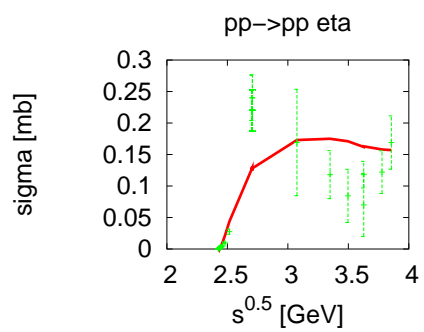
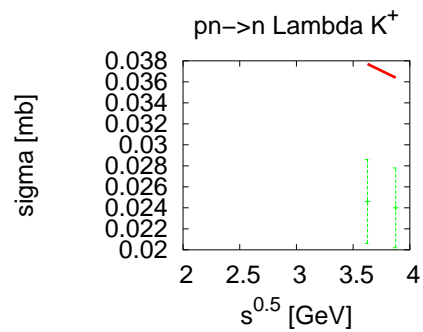
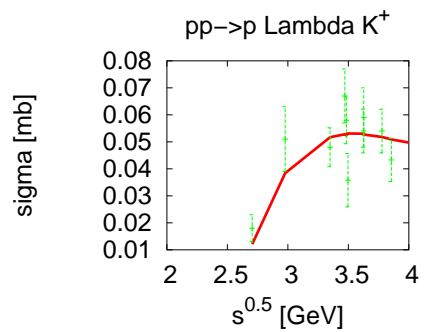


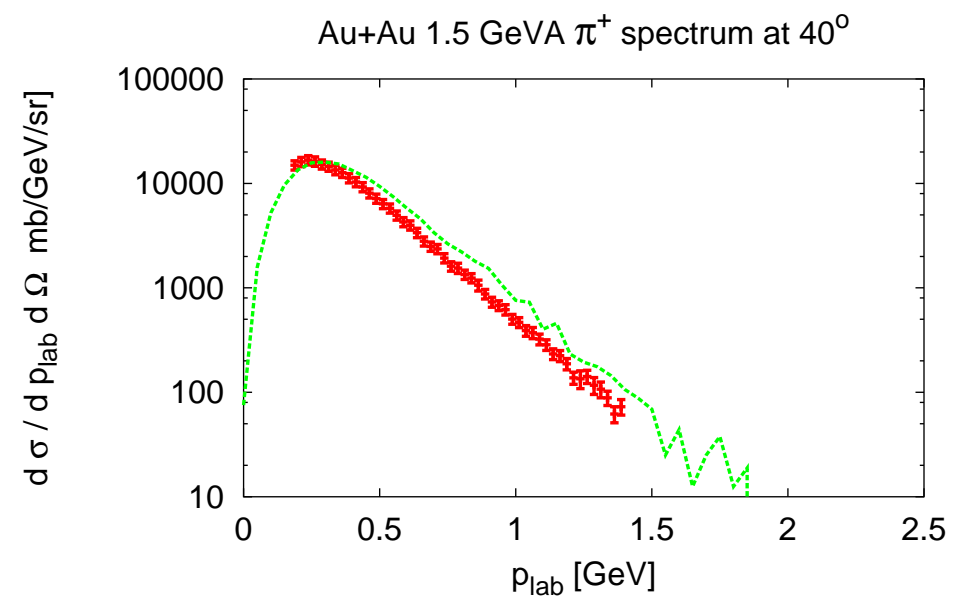
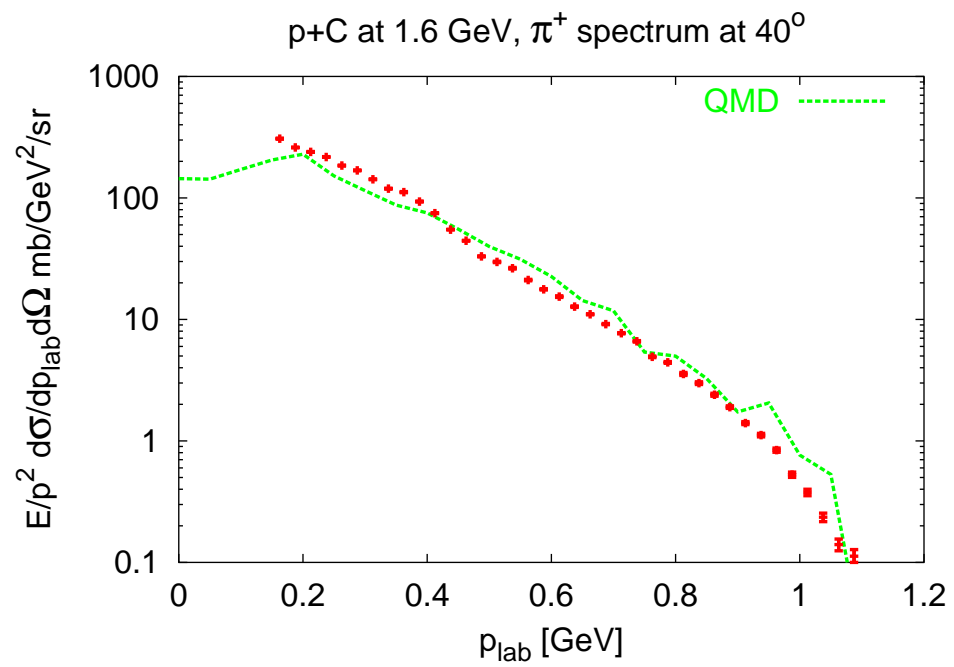
pp→pp omega



pp→pp pi⁺ pi⁻







ρ-A1 mixing

- $\overline{\Pi}_\rho = (1 - x)\Pi_\rho^v + x\Pi_{A1}^v$
- $\overline{\Pi}_{A1} = (1 - x)\Pi_{A1}^v + x\Pi_\rho^v$

where $x = \frac{1}{2}0.3 * \frac{\text{density}}{\text{normal density}}$

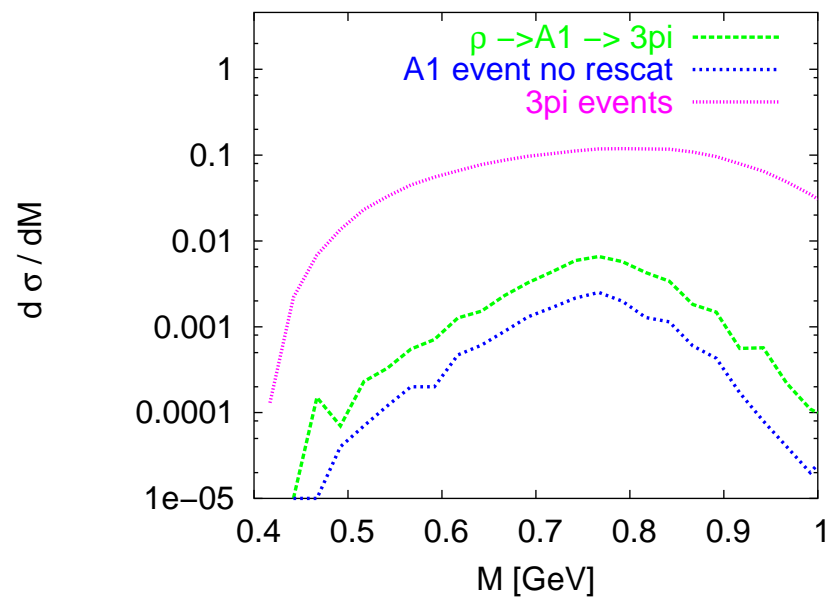
0.3 comes from chiral symmetry

- Signal: 3π decay of A_1 with the mass of ρ -meson
- π -nucleus collisions
- 1 million central event

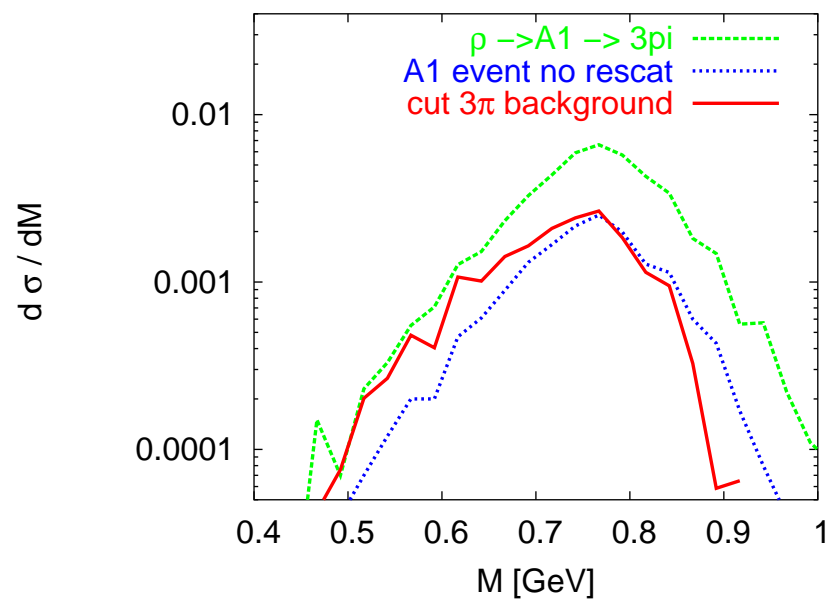
system	energy	A1	3 pions	A1 no rescatt.
$\pi^- Au$	1.3 GeV	6400	56 000	100
$\pi^- C$	1.3 GeV	1200	43 000	450
$\pi^- Ca$	1.3 GeV	3200	54 000	400
$\pi^- Ca$	1.1 GeV	3100	46 000	350
$\pi^- Ca$	1.5 GeV	2400	55 000	360

- Light system is preferred

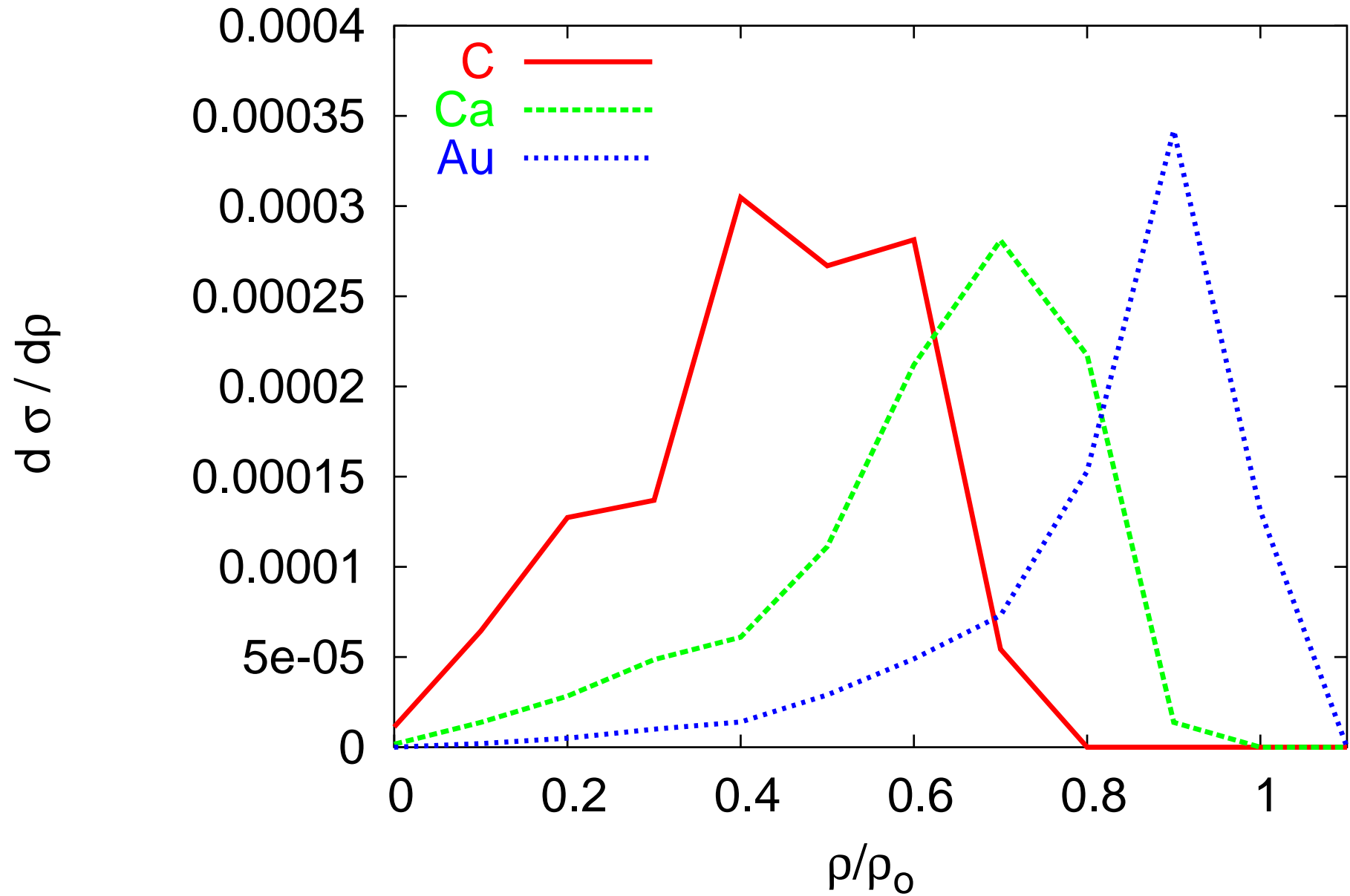
$\pi^- C$ 1.3 GeV A1 π^-



$\pi^- C$ 1.3 GeV A1 π^-



π^- A 1.3 GeV: A_1 dens. dep.



Summary

- The in-medium modification of the ρ is controversial
- A possible signal of chiral restoration in $A_1 - \rho$ mixing
- Needs very large statistics
- several theoretical uncertainties