



Új Nemzeti  
Kiválóság Program

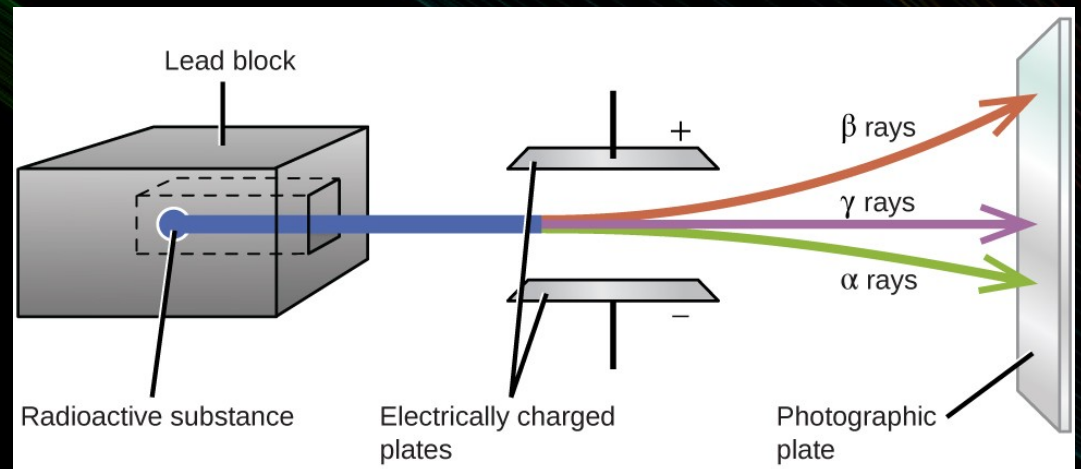
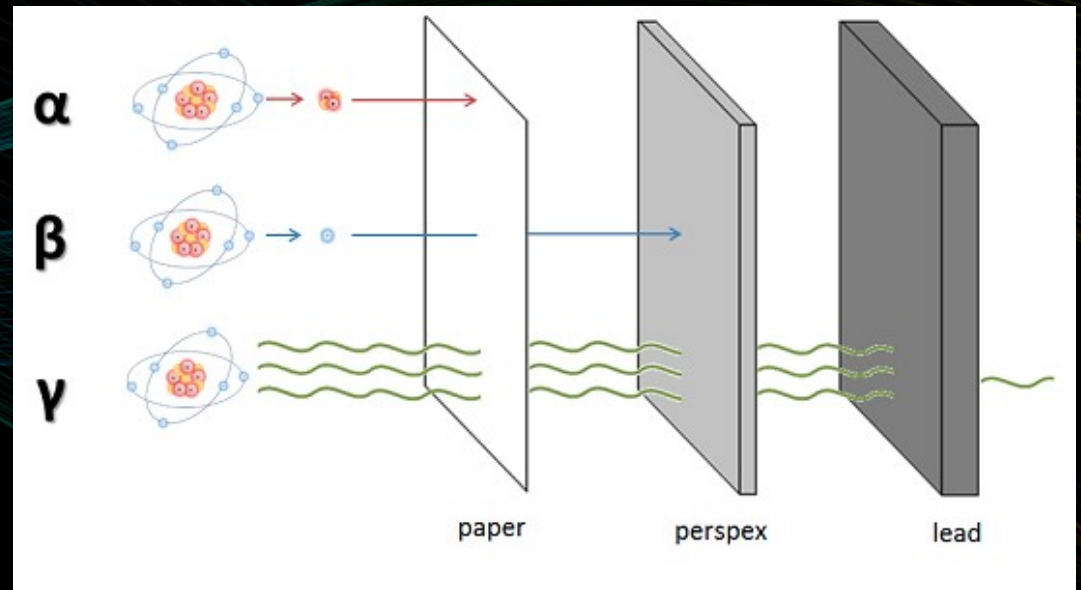
# IV. Radioactivity

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# Types of radioactivity

- $\alpha$ -radioactivity:  ${}^4\text{He}$  ions
  - short range (a paper absorbs)
  - high ionization power: dangerous in case of incorporation!
- $\beta$ -radioactivity:  $e^-$  or  $e^+$ 
  - midrange  $\rightarrow$  absorbed by an mm thick Al sheet (40-60 cm in air)
- $\gamma$ -radioactivity: photon
  - energetic electromagnetic radiation from the excited nucleus
  - very long range, low ionization power



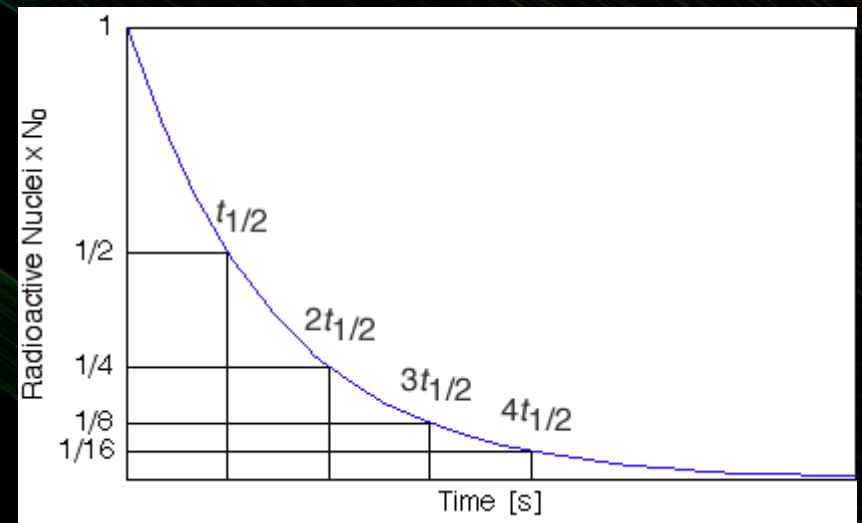


# Radioactivity: decay law

- A random (statistical) process with constant probability
- The number of decays proportional to the number of radioactive nuclei:

$$\frac{dN(t)}{dt} = -\lambda N(t) \longrightarrow N(t) = N_0 e^{-\lambda t}$$

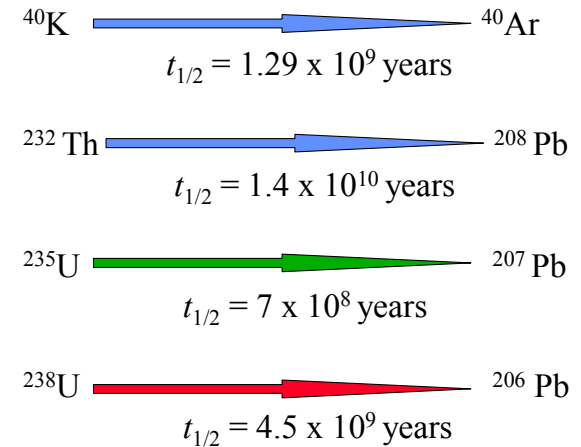
- $\lambda$  is the decay constant: „speed” of the decay
- Mean lifetime ( $\tau$ ): 1/e of nuclei undergo decay  $\rightarrow \tau = 1/\lambda$ 
  - level width:  $\Gamma = \hbar/\tau$
- Half life ( $T_{1/2}$ ): half of the nuclei undergo decay  $\rightarrow T_{1/2} = \ln 2/\lambda$
- Activity:  $a(t) = -\lambda N(t)$ 
  - unit is Becquerel (Bq): 1 decay / second



# Nuclear decay chains

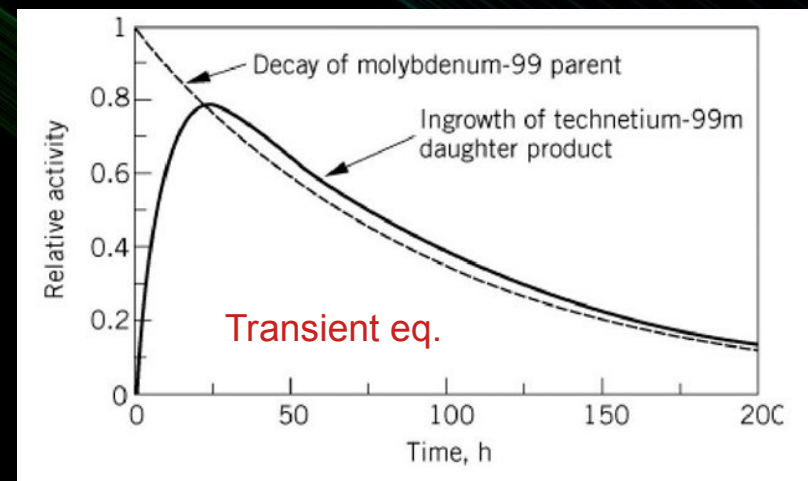
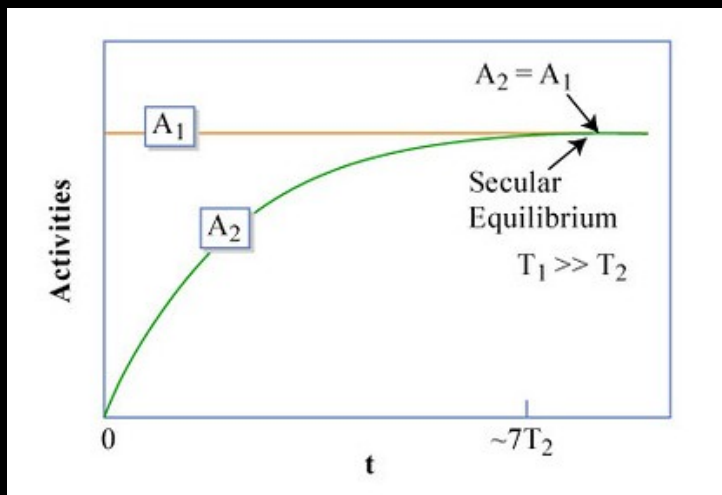
- Naturally occurring radioactive isotopes on Earth (and in our body even!)
- $^{40}\text{K}$  is beta-radioactive while the rest is mostly alpha active
- The final products are Pb isotopes ( $Z=82$ , shell structure again...)
- Activity of different members within the decay chain can be calculated with differential equations:

## Natural Decay Series of Existing Isotopes



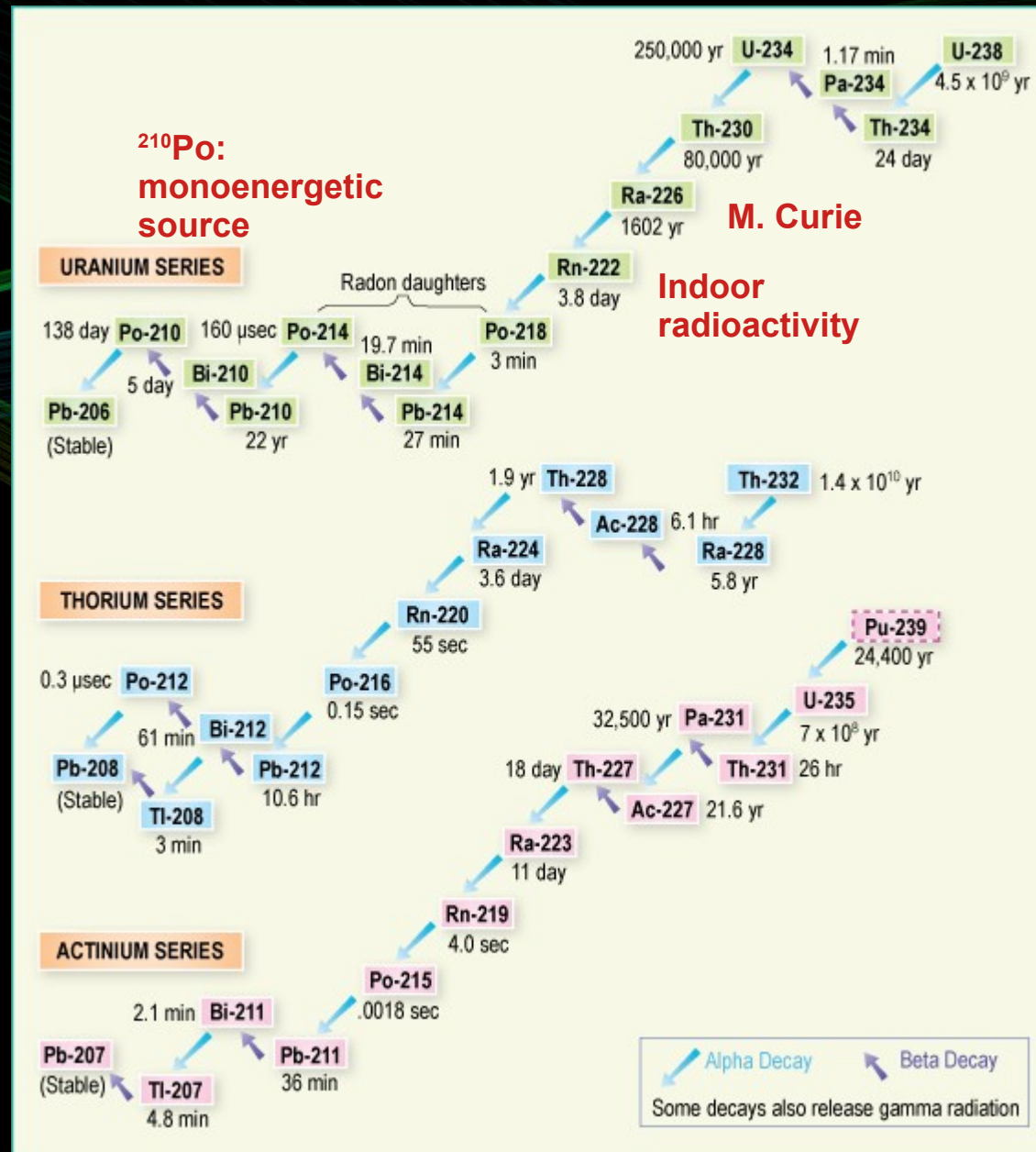
$$\frac{dN_1(t)}{dt} = -\lambda N_1(t)$$

$$\frac{dN_2(t)}{dt} = \lambda_1 N_1(t) - \lambda_2 N_2(t)$$



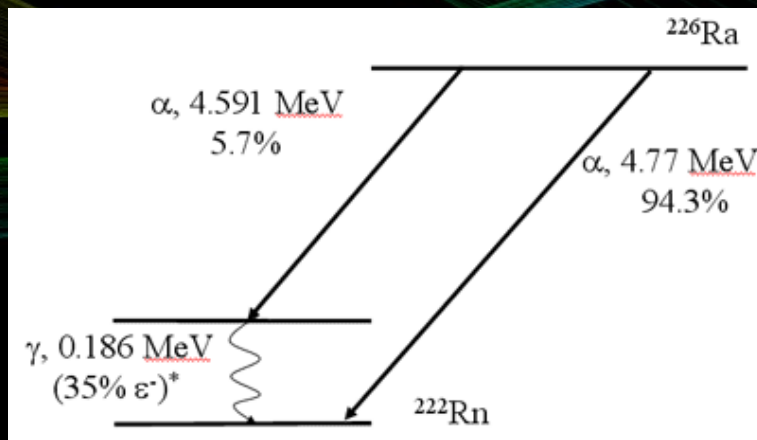
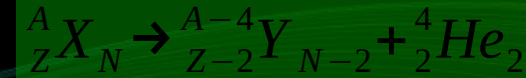


# Nuclear decay chains



# Alpha decay

- $\alpha$ -particle is a  ${}^4\text{He}$  nucleus ( ${}^4\text{He}^{++}$  ion)
- the energy balance of alpha decay and the condition for alpha decay:  
 $Q = M_x(A, Z) - [M_y(A-4, Z-2) + M_{\text{He}}(4, 2)] > 0$  (from LDM  $\rightarrow Z > 73$ )

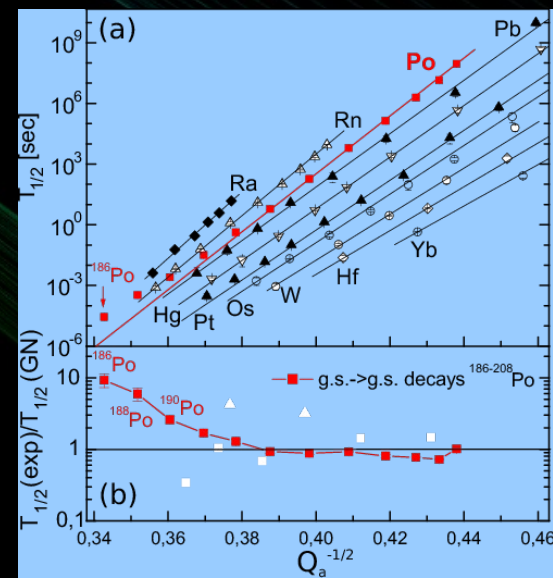


- $\alpha$ -decay of  ${}^{226}\text{Ra}$
- $\alpha$ -particle energy is determined by the initial and final states
- transitions to higher (excited) levels are possible (mostly to low lying rotational states of deformed daughter nuclei) but with quickly decreasing probability

- $4 \text{ MeV} < E_\alpha < 9 \text{ MeV}$  while  $10^{-7} \text{ s} < T_{1/2} < 10^{10} \text{ year}$  (!!)
- relation between  $E_\alpha$  and  $T_{1/2}$  within the 3 natural decay chains is the Geiger-Nuttall rule (1911)

$$\lg \lambda = a \cdot \lg E_\alpha + b$$

where  $a$  is the same for all chains,  $b$  is not





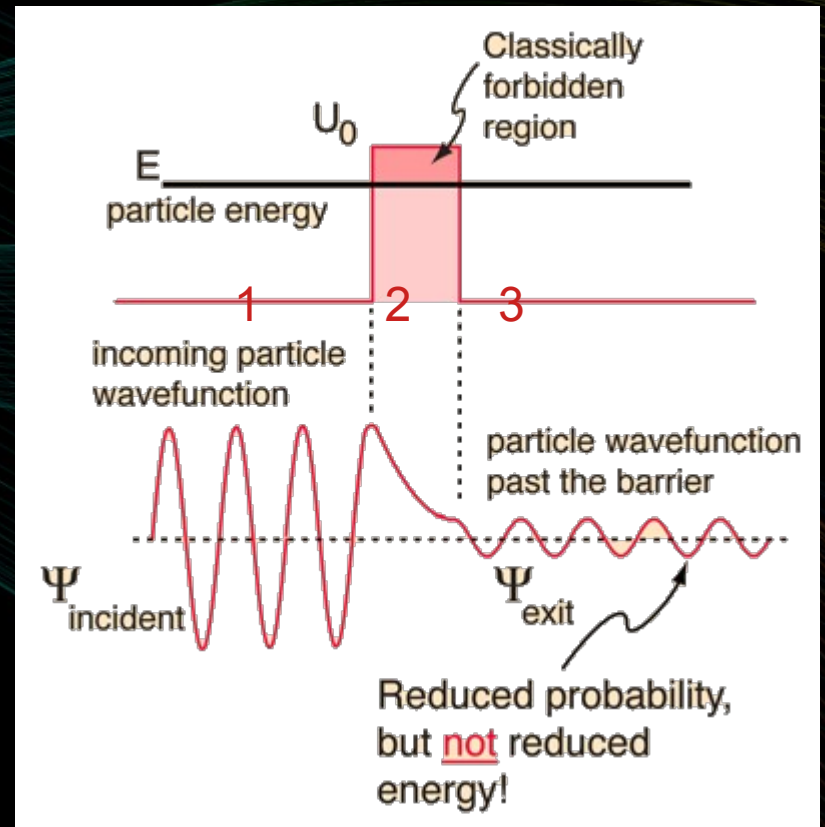
# Alpha decay – quantum tunnelling

- *Gamow*: understanding the Geiger-Nuttall rule within quantummechanics
- From the solving the Schrödinger equations in the 3 different region (where  $U_0 < E$  [1-3] and  $U_0 > E$  [2])

$$\Delta \Psi + (2m/\hbar^2)(E - V)\Psi = 0$$

$$D = \exp\left[-\frac{2}{\hbar} \sqrt{2m(U_0 - E)} x_0\right]$$

transition probability



- If  $U_0 = Z_{\text{nucleus}} z_{\alpha} e^2/r$  is the Coulomb potential (3 dimensional) and  $E_{\alpha}$  is the energy of the  $\alpha$ -particle  $\rightarrow$  GN rule can be understood

- If  $l_{\alpha} > 0$  (if it is allowed, see next)  $\rightarrow$  probability is decreased by the centrifugal barrier:

$$V_{cf}(r) \approx \frac{l(l+1)\hbar^2}{3mr^2}$$

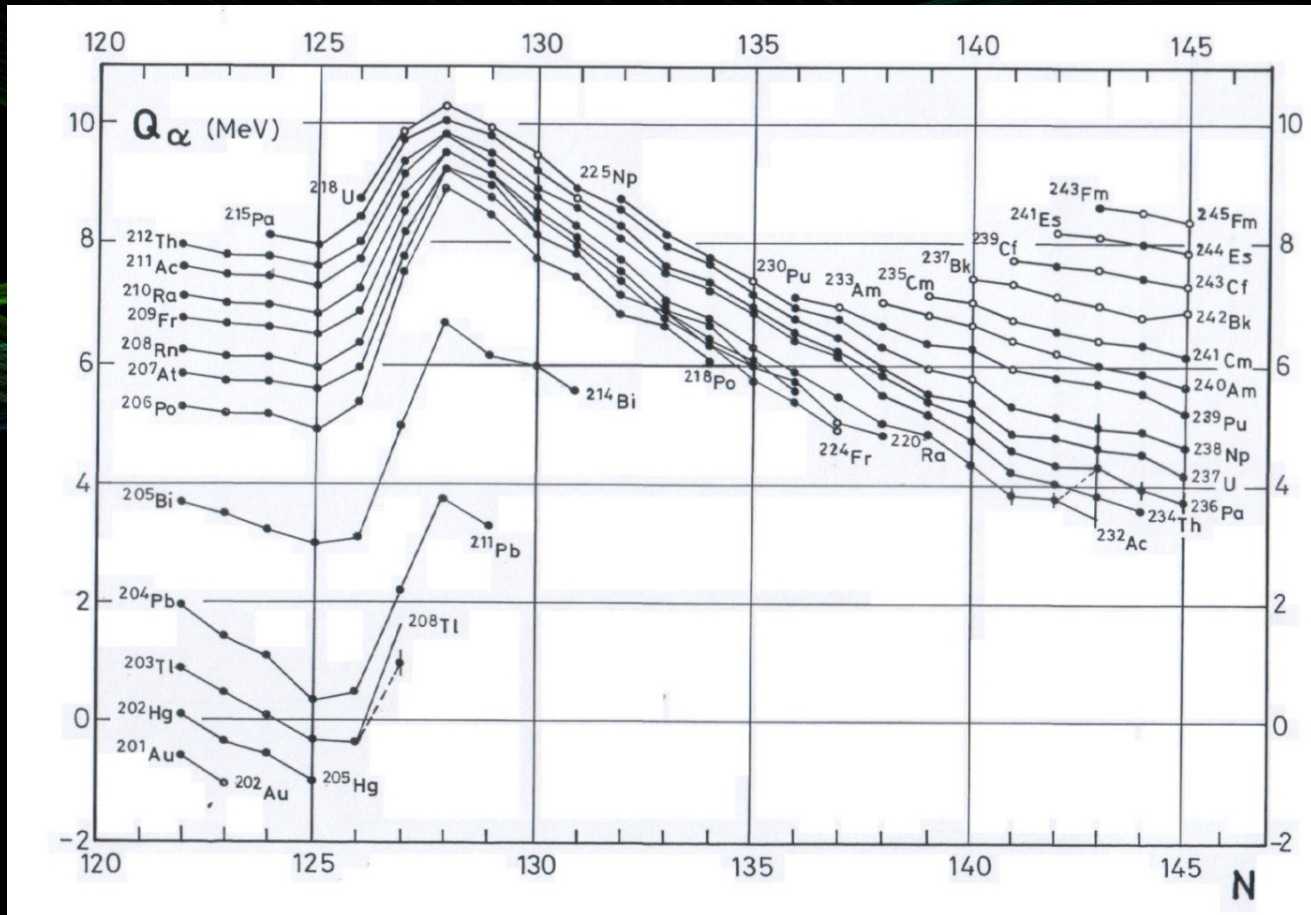
# Alpha decay

- Controlled by the strong force
- Kinematics: momentum conservation
  - $|p_\alpha| = |p_{\text{daughter}}| \rightarrow E_\alpha = Q_\alpha [M_{\text{daughter}} / (M_{\text{daughter}} + M_\alpha)]$  ( $\rightarrow E_{\text{daughter}}$  is typically small:  $\sim 2\%$ )
- **P** parity and **T** isospin is conserved  $\rightarrow$  selection rules
  - $T_\alpha = 0 \rightarrow$  parent and daughter nuclei have the same isospin
  - $P_\alpha = +1$  and  $I_\alpha = 0 \rightarrow ||I_p - I_d| \leq I_\alpha \leq |I_p + I_d|$  and  $P_p = (-1)^{I_\alpha} P_d$
- Some remarks:
  - From D to  $\lambda$  :  $\lambda = P\nu D$  (P: probability of an  $\alpha$  formation within the nucleus and  $\nu$ : the frequency of the  $\alpha$  particle „hitting” the surface)  $\rightarrow$ 

$$\lg \lambda = \frac{A}{\sqrt{E_\alpha}} + B$$
  - A forbidden alpha transition is not „absolutely” forbidden (but decreased in probability by a factor of  $10^7$ - $10^{14}$ )
  - Centrifugal barrier is not so large  $\rightarrow$  only little dependence on  $I_\alpha$



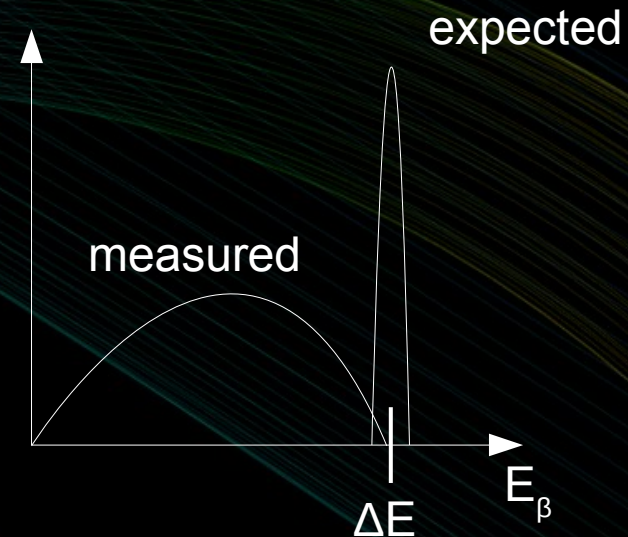
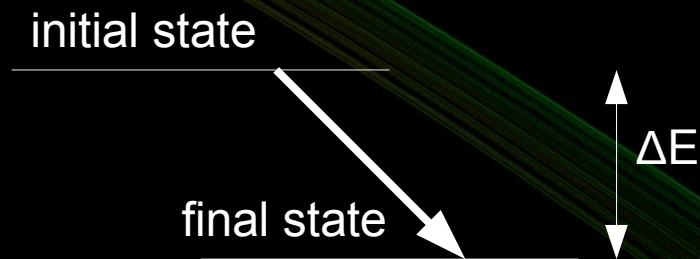
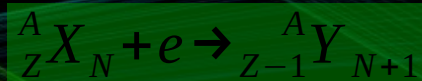
# Alpha decay systematics



- Large  $Q$  at  $N=128 \rightarrow$  shell structure

# Beta decay

- Transmutation of an unstable nucleus to an isobar one with  $\Delta Z = \pm 1$
- Condition:  $E_{\beta^-} = [M_{\text{at}}(A, Z) - M_{\text{at}}(A, Z+1)]c^2$  and  $E_{\beta^+} = [M_{\text{at}}(A, Z) - M_{\text{at}}(A, Z-1) - 2m_e]c^2$
- Experimental observations:  $0.02 < E_{\beta} < 16 \text{ MeV}$  and  $10^{-2} \text{ s} < T_{1/2} < 10^{15} \text{ years}$

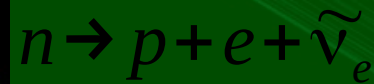


- Experimental technique: magnetic spectrometers
- Instead of  $E_{\beta} = \Delta E$  a continuous  $E_{\beta}$  is measured with  $0 < E_{\beta} < \Delta E$ !
- maximum yield is at  $E_{\text{max}}/3$  (for natural radioactivity  $E_{\beta} = 0.25 - 0.45 \text{ MeV}$ )
- more symmetric  $\beta$ -spectrum for light nuclei:  $E_{\text{max}}/2$



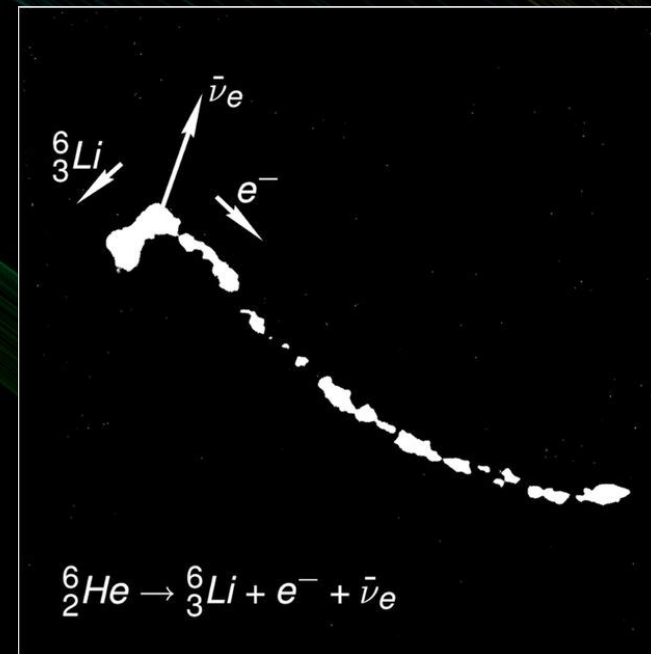
# Beta decay

- Explaining the continuous  $E_\beta$  spectrum:
  - monochromatic  $e^-$  are emitted, but scattering on the electron „clouds“ of atoms
  - no energy conservation in  $\beta$  decay....
  - W. Pauli (1931): a third particle takes away the missing energy!
    - the particle should have zero charge (charge conservation), small mass and  $s=1/2\hbar$



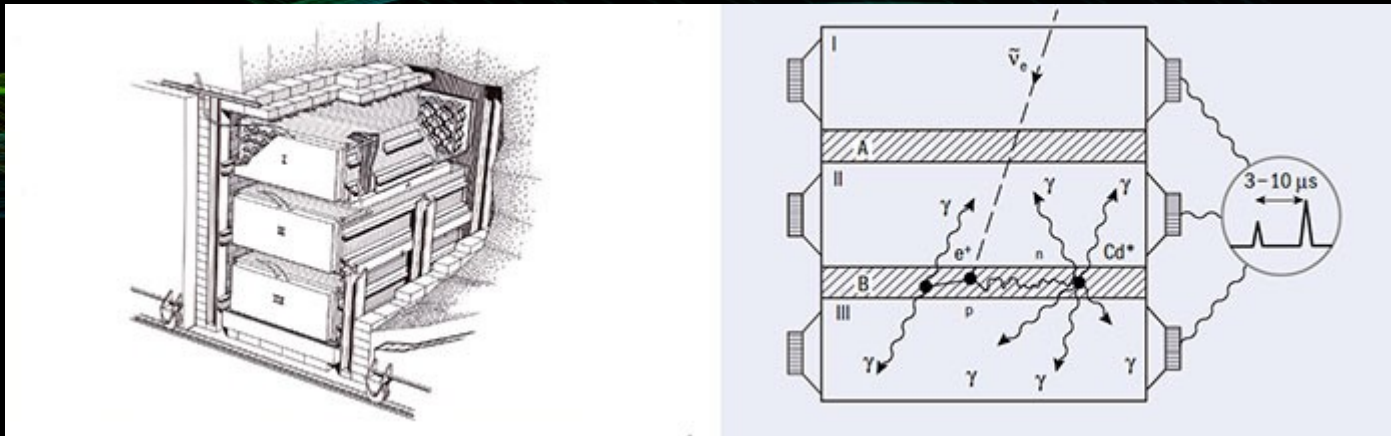
neutrino hypotheses

- Szalay - Csikai experiment (1956):
  - indirect observation of neutrino
  - momentum conservation should be satisfied
  - photon of a decay in a Wilson cloud chamber in Debrecen
  - (unfortunately published only in 1957)



# Beta decay

- *Reines and Cowan* (1956): inverz beta decay of neutrons

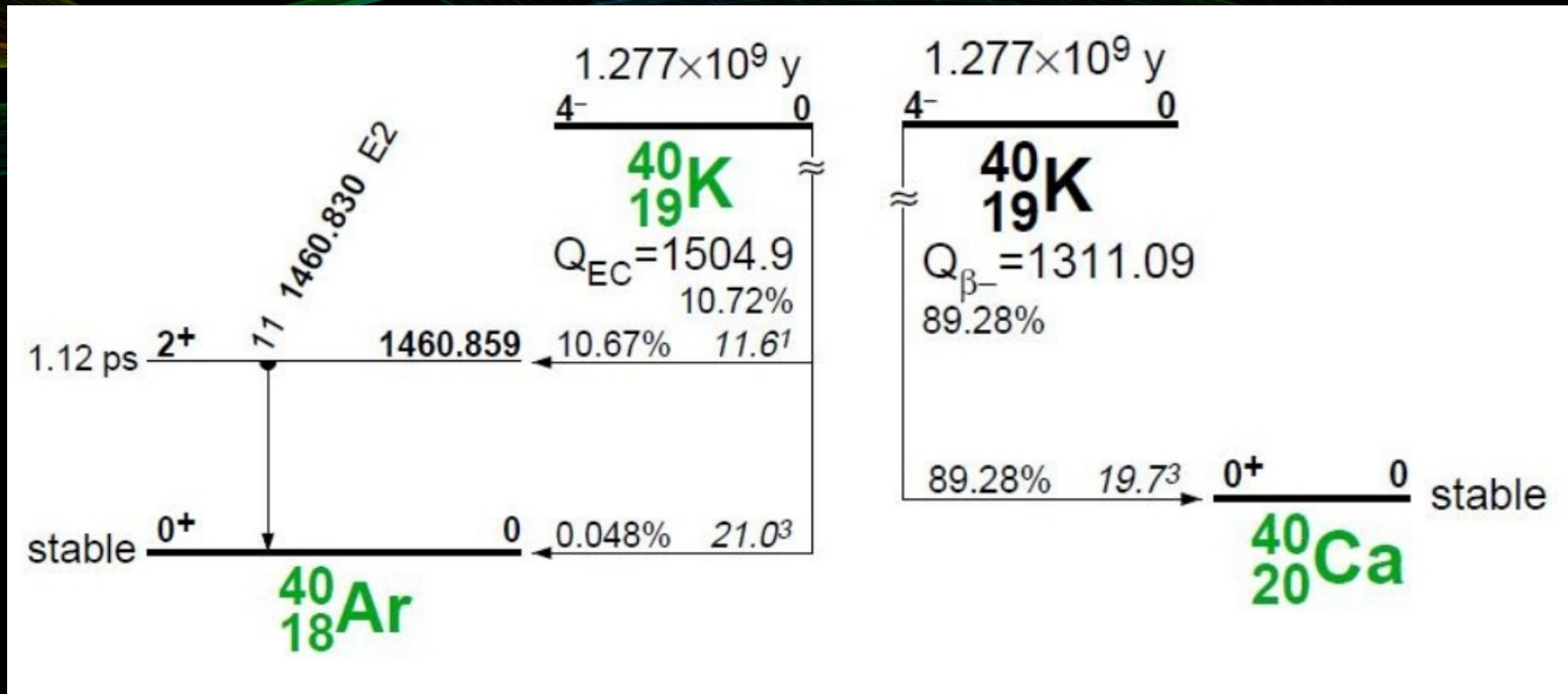


- *A, B* volumes:  $CdCl_2$  water-solution  $\rightarrow$  Cd neutronabsorber, water moderator
- *I* and *II* and *III* volumes: liquid scintillator material (1 m<sup>3</sup>)
- Triple coincidences of
  - the positron annihilation (2x511 keV  $\gamma$  photon)
  - neutrons, after moderation, are absorbed by Cd which gives delayed  $\gamma$  photons
- Cross section was determined to be  $\sigma=10^{-43}$  cm<sup>2</sup>  $\rightarrow$  10<sup>17</sup> km mean free path in condensed material...



# Beta decay: an example

- decay scheme of  $^{40}\text{K}$ : both type of beta decay!
- transitions to higher levels  $\rightarrow$  gamma decay



# Beta decay

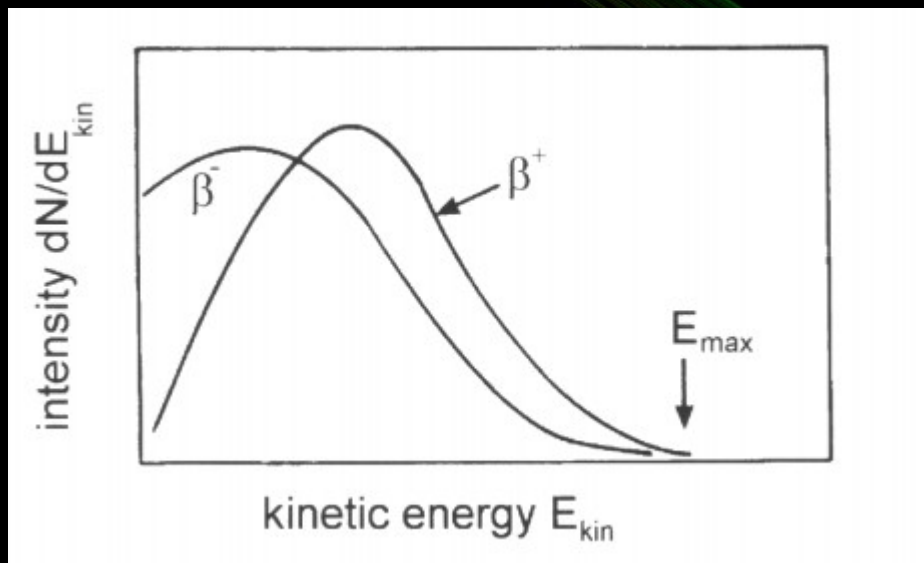
- Fermi (1934): understanding the beta spectrum
  - beta decay lifetime is typically long → the interaction is weak → perturbation theory can give the transition probability
  - electron and neutrino can be in all the quantum states with equal probability
  - the probability of emitting an electron with  $E - E+dE$  energy:

$$P(E) dE = 2 \frac{\pi}{\hbar} |H_{i,f}|^2 \rho(E) dE$$

$\rho(E)$ : final state density  
 $H_{i,f}$ : transition matrices (depend on the initial, the final state and the interaction)

$$\rho(E) = \sqrt{E^2 - m^2 c^4} \sqrt{(E_0 - E)^2 - m_\nu^2 c^4} E (E_0 - E)$$

(if  $m_\nu = 0$  than simpler)



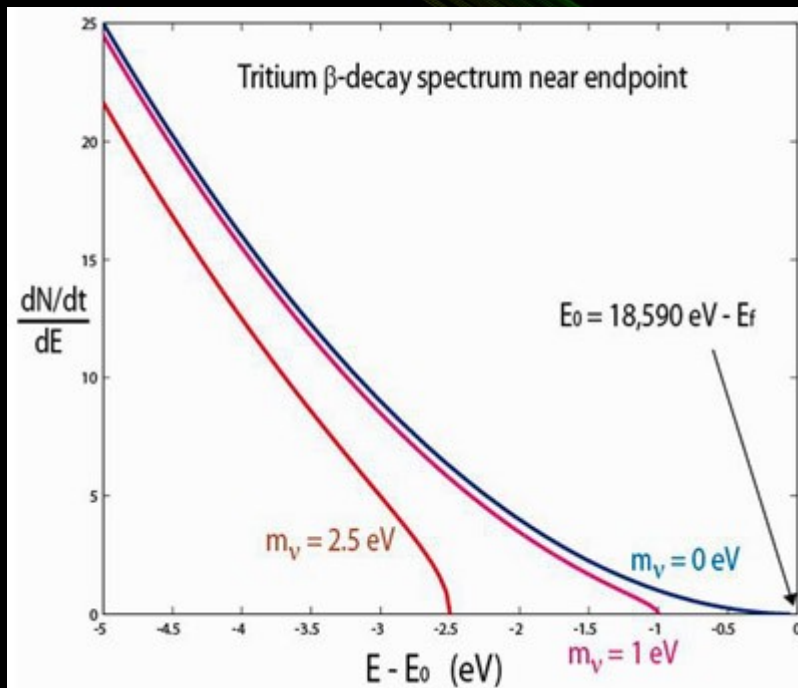
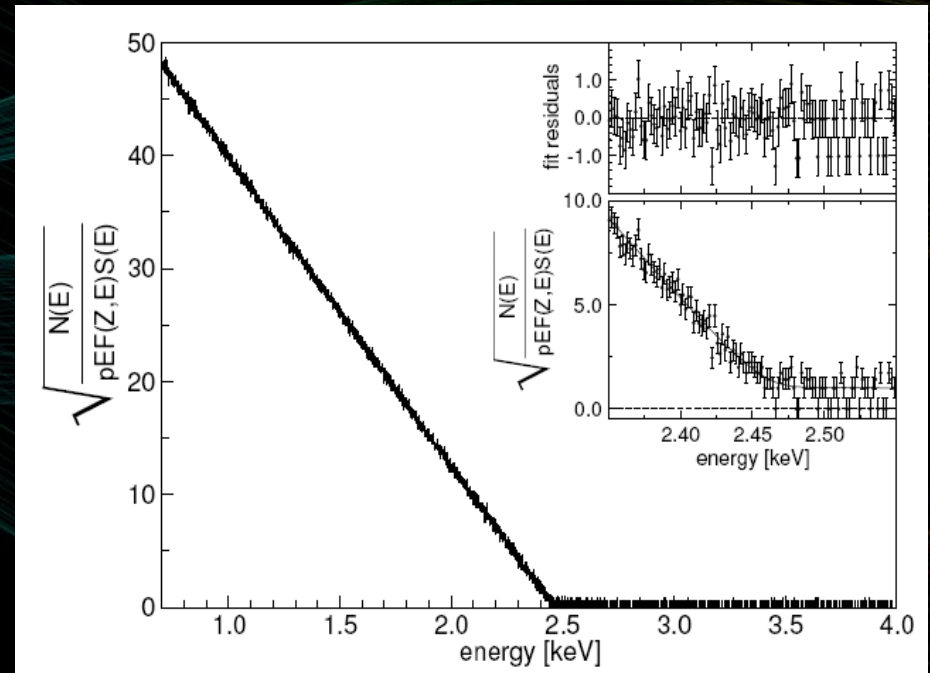
- Spectrum starts at  $mc^2$  (E is the total energy above)
- Coulomb correction: nucleus is positively charged → attracts electrons, pushes positrons → spectrum is shifted / distorted
- Fermi introduced a factor:  $F(Z,E)$



# Beta decay theory

- Fermi-Kurie plot: for convenience it is better the plot a normalized form and not the spectrum  $N(E)$

$$\sqrt{\frac{N(E)}{E \sqrt{E^2 - m^2 c^4} F(Z, E)}}$$



- for allowed beta decay, it is linear
- $E_0$  is defined by the crossing with the X-axis  $E_e$
- if neutrino has mass, than the electron energy end-point differs from  $E_0$  by  $m_\nu c^2$
- see the scale!!

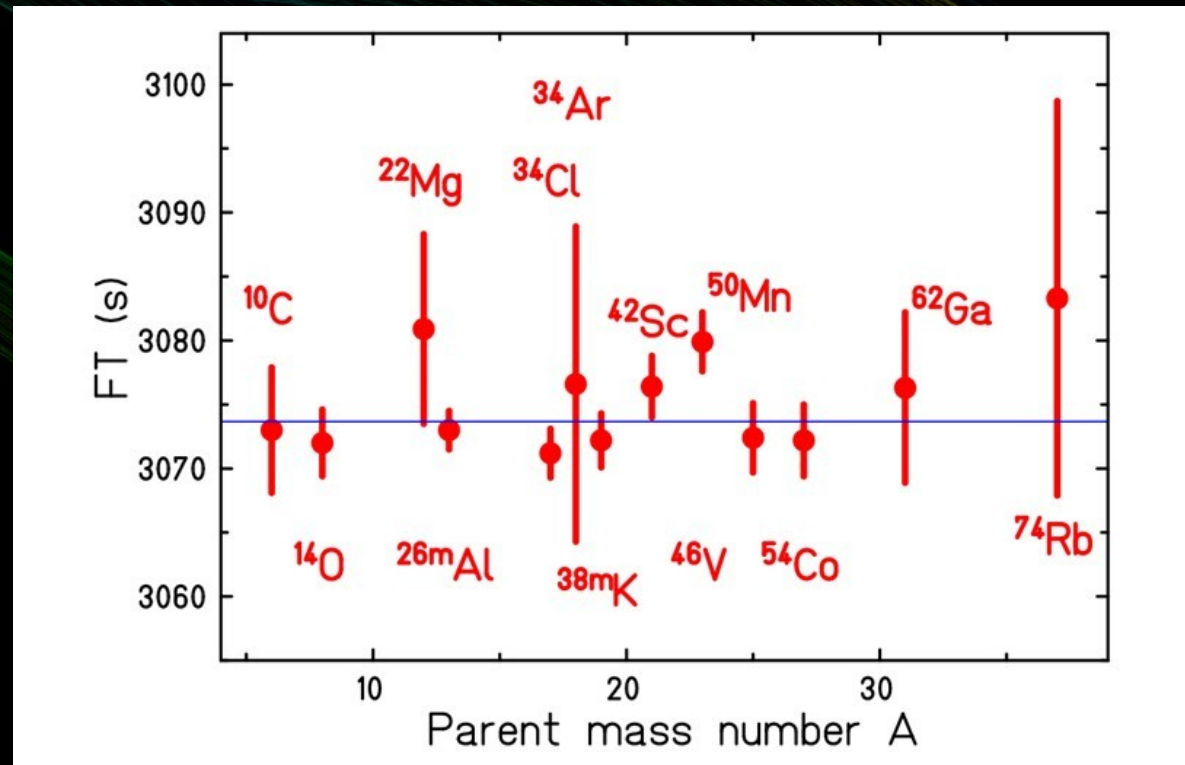
# Beta decay: theory

- Total probability of beta decay:

$$\lambda = \frac{1}{\tau} = f(Z, E_0) = \int_0^{E_0} \rho(E) F(Z, E) dE \approx \int_0^{E_0} E^4 dE = E_0^5 \quad \Rightarrow \quad f \cdot \tau \approx \text{constant}$$

- $f(E_0)\tau \rightarrow$  comparative half life is constant for superallowed, allowed, and forbidden transitions

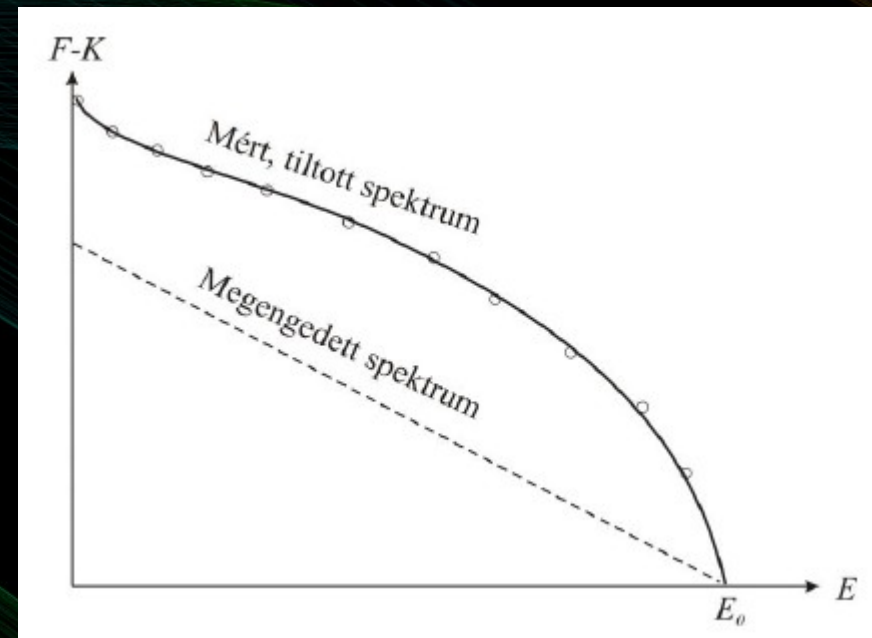
transition type	log(ft)
superallowed	2.9-3.7
allowed	4.4-6.0
first forbidden	6.0-10
second forbidden	10-13





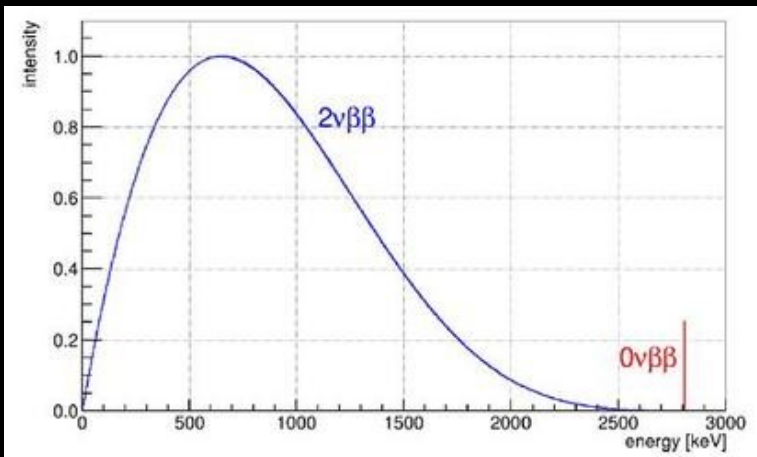
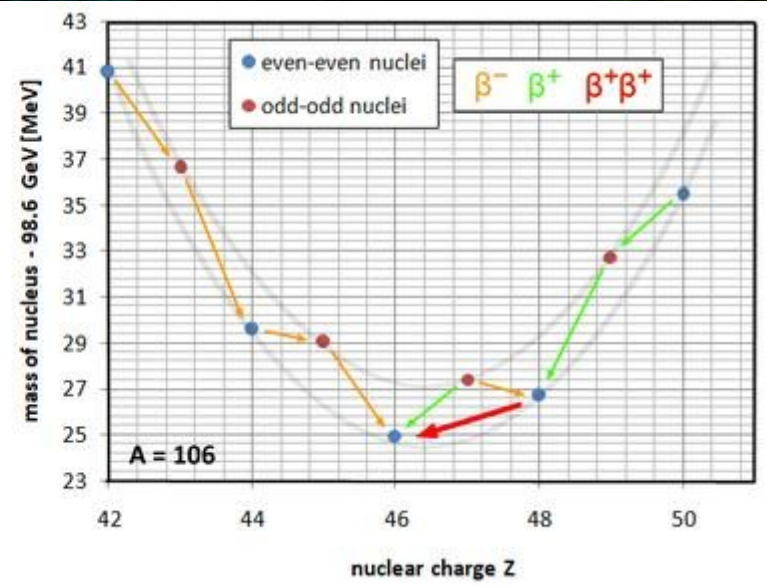
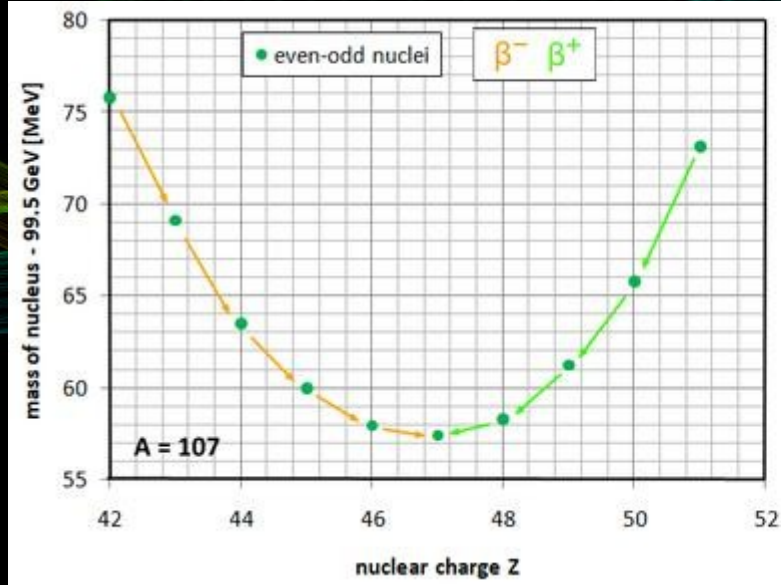
# Beta decay: theory

- Allowed transitions:  $l=0$  and  $P_1=P_2$ 
  - Fermi transition:  $s_e+s_v = 0$  (spins are antiparallel)
  - Gamow-Teller transition:  $s_e+s_v = 1$  (spins are parallel)
    - Sun:  $2p \rightarrow d + e^+ + \nu$
    - Pauli principle:  $s(2p)=0$  but  $s(d)=1$  !!
    - Fermi transition is not possible.
- Forbidden transitions:  $l > 0$  and/or  $P_1 \neq P_2$ 
  - transition probability is very small
  - Kurie plot is deviated from linear
  - characterized numerically by the  $\lg(Ft)$  value



# Double beta decay

- Sometimes energetically beta decay is not allowed between even-even isobars → but it is for Z+2!! → double beta decay is allowed



(if  $\bar{\nu} \neq \nu$ )  $2\nu\beta\beta: 2n \rightarrow 2p+2e^{-}+2\bar{\nu}$

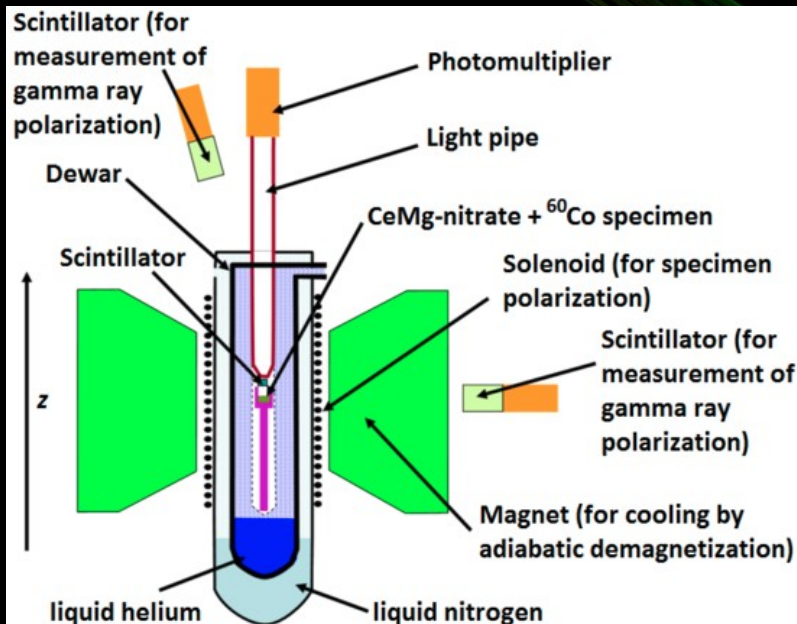
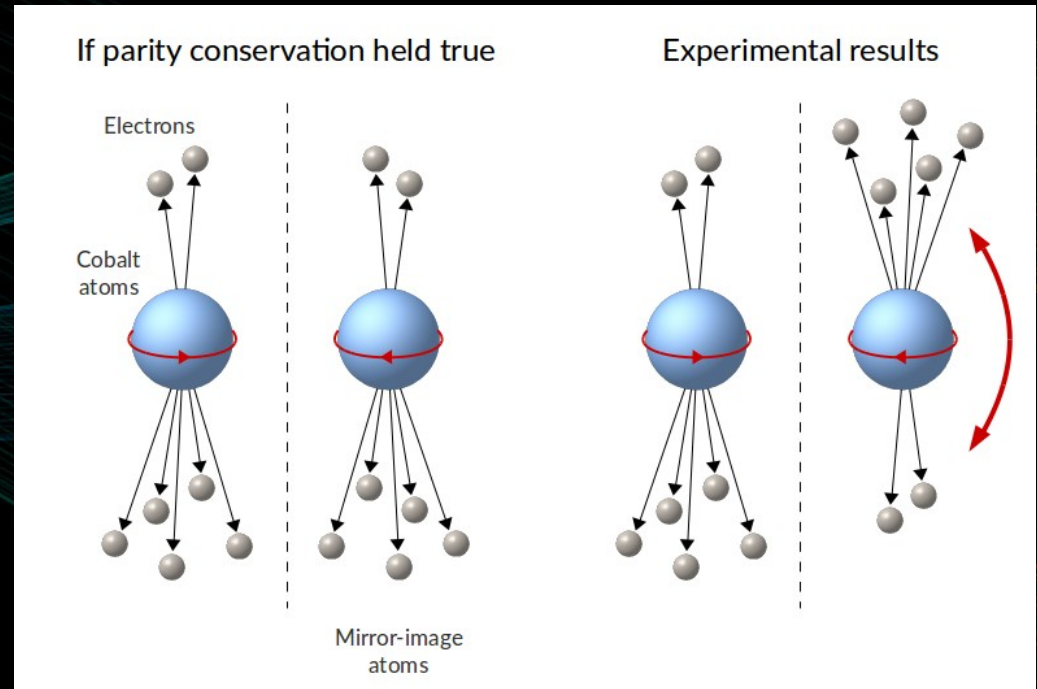
(if  $\bar{\nu} = \nu$ )  $0\nu\beta\beta: n \rightarrow p+e^{-}+\bar{\nu} + n \rightarrow p+e^{-}$

- N(E) gives the transition energy E0
- Neutrino is a majorana particle? Very intense research field today!



# Beta decay: parity violation

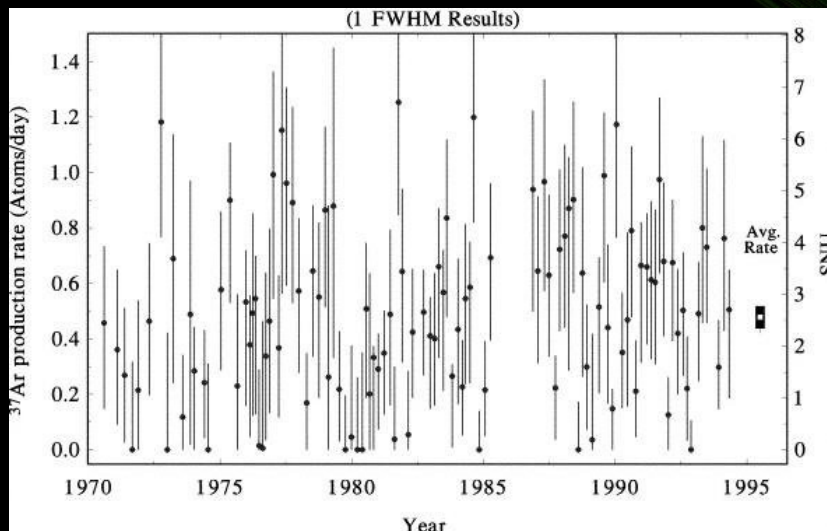
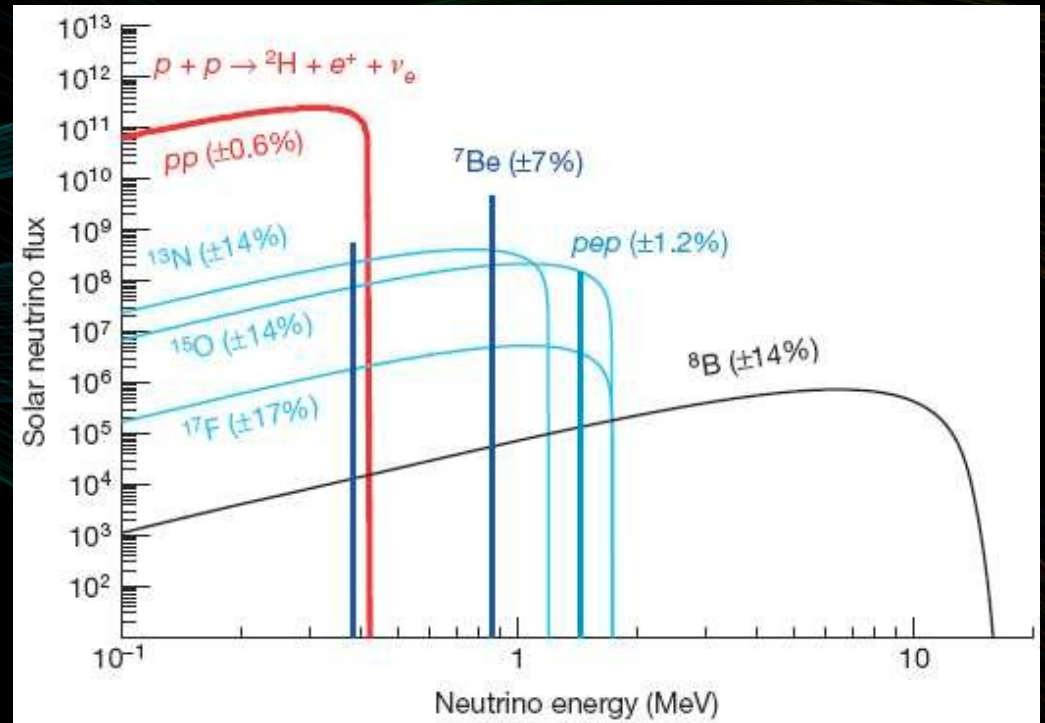
- C. Wu experiment (1956):
  - in strong and EM interaction parity is conserved
  - what about weak interaction, in beta decay?



- Top science, top technology!
  - needed high magnetic field (10T), and very low temperature to align  ${}^{60}\text{Co}$  spin
  - $T=0.003\text{K}$  (!! ) was achieved by adiabatic cooling
- Result:  $e^-$  emission preferred the direction which is opposite to the  ${}^{60}\text{Co}$  spin!!

# Neutrino physics

- R. Davis (1969): measured the solar neutrino flux
- Sun is an extreme source of neutrinos: 64 billions of neutrinos per second per  $1\text{cm}^2$  !!
  - But very low cross section:  $10^{17}$  m
- pp-cycle: 98.5% of power (rest is CNO cycle)
  - almost all neutrinos are low energy  $E < 0.42$  MeV
  - but  $^8\text{B}$  and  $^7\text{Be}$  neutrinos are energetic:  $E < 14$  MeV

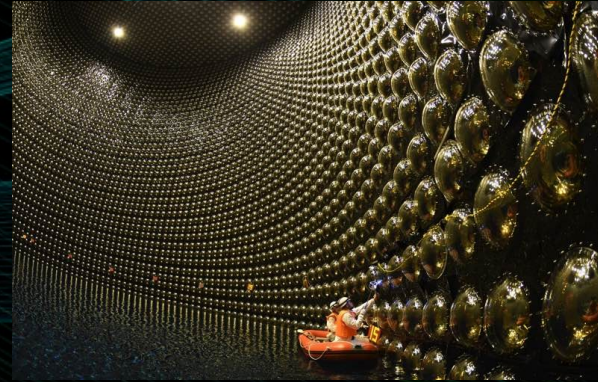


- The experiment was at Homestake mine with 615 tons of tetrachlore-etan,  $\text{C}_2\text{Cl}_4$  a cleaning fluid
  - $^{37}\text{Cl} + \nu_e \rightarrow ^{37}\text{Ar} + e^- \quad E > 0.814 \text{ MeV}$
- $^{37}\text{Ar}$  decay was detected
- Measured 0.48 atoms/day instead of 1.5 atoms/day

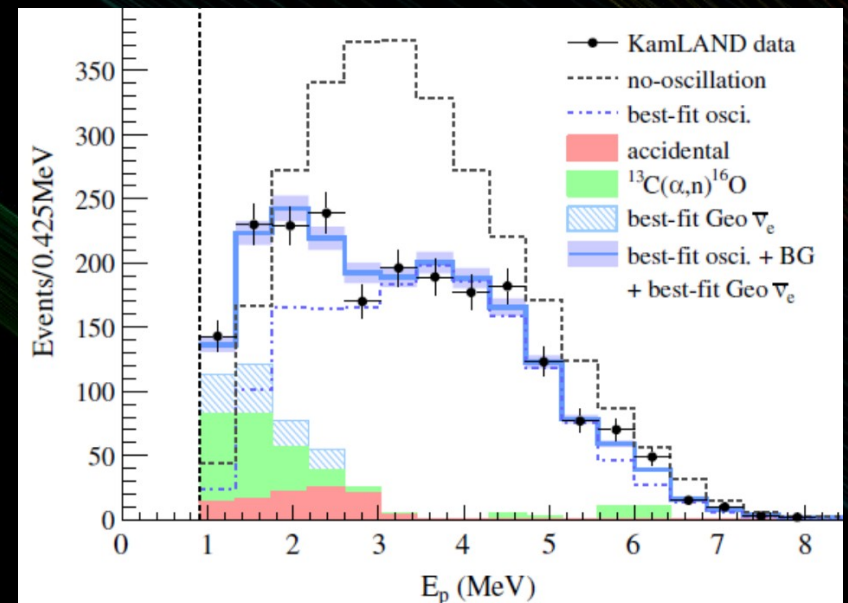


# Neutrino physics

- Further experiments:
  - Solar and atmospheric: Super-Kamiokande, Japan, 1998
  - Solar neutrinos: Sudbury Neutrino Observatory, Canada, 2001
    - sensitive to all type of neutrinos

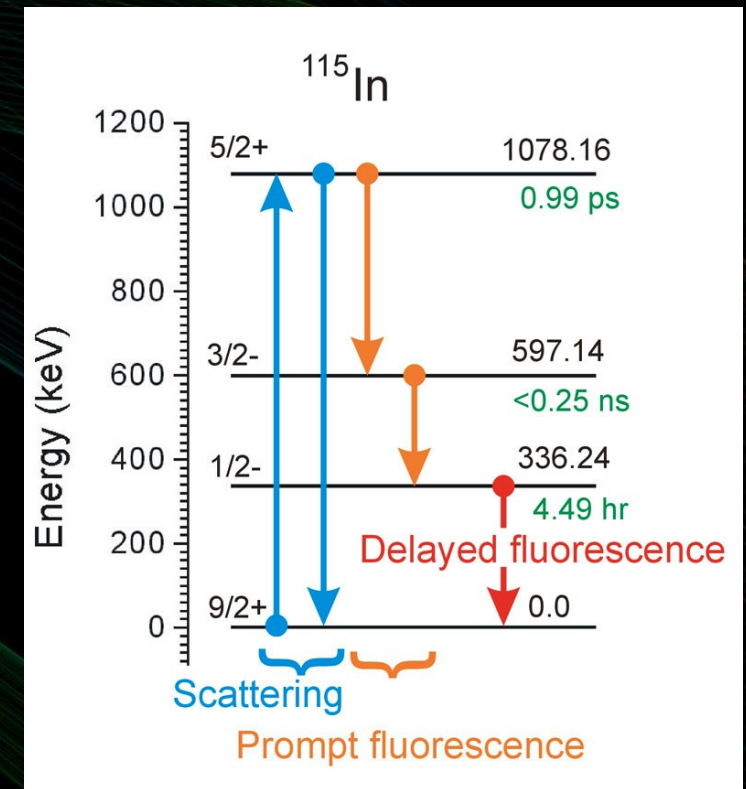
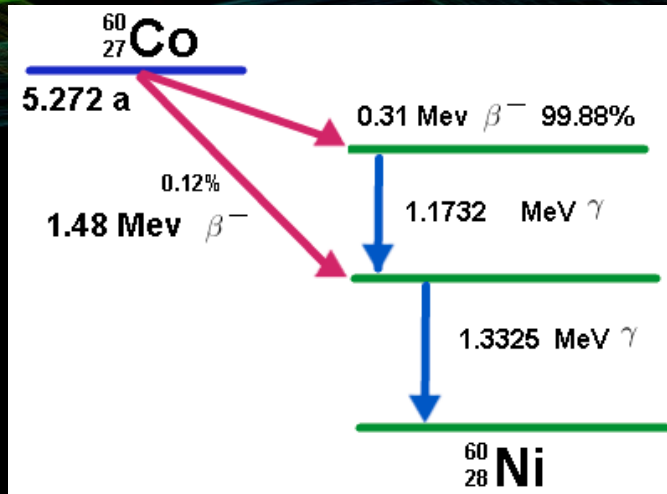


- Solution: neutrino oscillations
  - electron, muon and tau neutrinos can transform to each other
  - detectors are sensitive only to electron neutrinos → the rest is „missing“
  - But! Oscillation only can take place if at least one of the neutrinos has mass!
- KamLAND experiment, 2005-2008



# $\gamma$ decay

- Excited states of a nucleus  $\rightarrow$  de-excitation by  $\gamma$  photon emission
- Electromagnetic radiation with short wavelength ( $10^{-8} > \lambda > 10^{-11}$  cm)
- Excitation is due to a)  $\alpha$  (typically very low lying excited states with  $E_\gamma < 0.5$  MeV  $\rightarrow$  quantum tunneling properties) b)  $\beta$ -decay to excited states or b) nuclear reactions (high energy excitations)



- One photon de-excitation  $\leftrightarrow$   $\gamma$  cascades
- $10 \text{ keV} < E_\gamma < 5 \text{ MeV}$  and  $10^{-16} < T_{1/2} < 10^{-8}$
- some exceptions: e.g.  $^7\text{Li} + p \rightarrow ^8\text{Be} + \gamma$ 
  - $E_\gamma = 17 \text{ MeV}$  (nucleon emission is forbidden by parity / angular momentum conservation)



# gamma decay

- Energy and momentum conservation:  $\mathbf{p}_\gamma + \mathbf{p}_M = 0$  and  $E_1 - E_2 = \Delta E = E_\gamma + E_M$
- Recoil energy of the nucleus:

$$E_N = \frac{E_\gamma^2}{2M_N c^2} \ll E_\gamma$$

(typically  $\sim 0.1 - 10$  eV)



Negligible. But not in Mössbauer spectroscopy!

- angular momentum conservation: selection rules

$$|I_1 - I_2| \equiv \Delta I \leq l \leq I_1 + I_2$$

- $l=1$  dipole;  $l=2$  quadrupole;  $l=3$  octupole radiation (no monopole radiation, photon spin is  $1\hbar$ )  $\rightarrow$  multipolarity is from angular distribution
- polarisation: electric (E) and magnetic (M)
- parity conservation:  $P_1/P_2 = (-1)^l$  (electric transition) and  $P_1/P_2 = (-1)^{l+1}$  (magnetic)

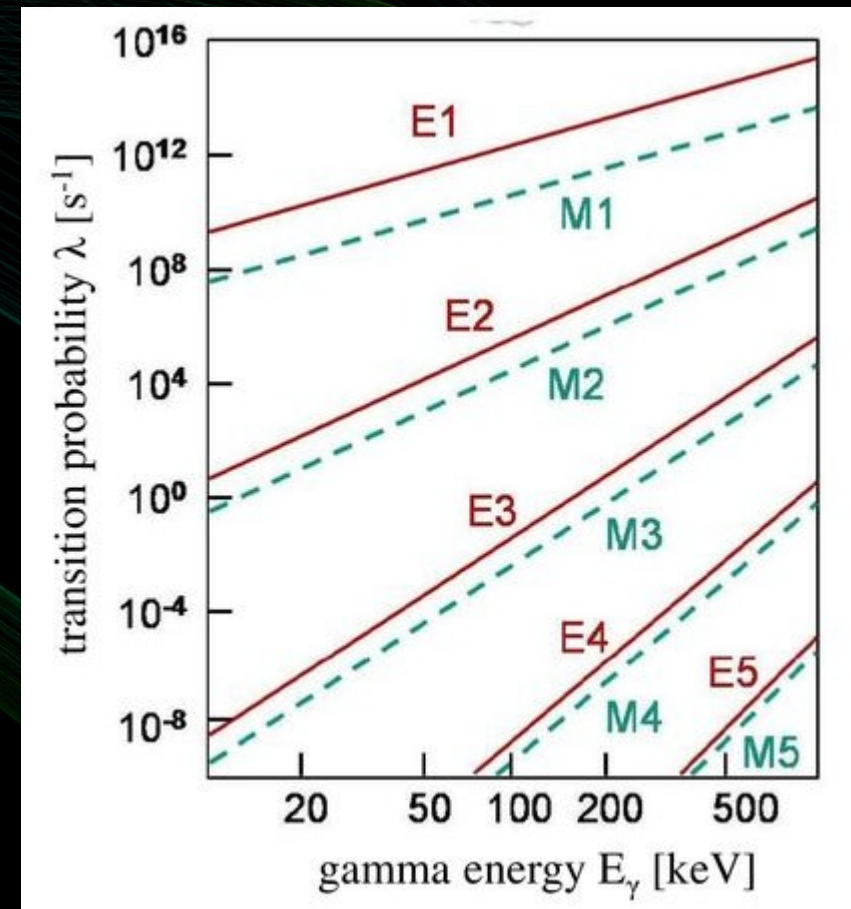
# gamma decay

- $\gamma$  decay probability strongly depends on the  $l$  multipolarity and  $E_\gamma$  transition energy:
  - increasing  $l \rightarrow$  decreasing transition probability (by a factor of  $10^{-4}$ )
    - dipole and quadrupole radiations dominate
    - large spin-difference in nucleus  $\rightarrow$  gamma decay probability is small (isomeric state)
    - transitions in a rotational band: cascades!
  - for given  $l$  magnetic transition probability is  $10^2$ - $10^3$  smaller than electric
- Weisskopf estimates: relation between  $E_\gamma$  and  $T_{1/2}$  for different multipolarities

$$|I_1 - I_2| \equiv \Delta I \leq l \leq I_1 + I_2$$

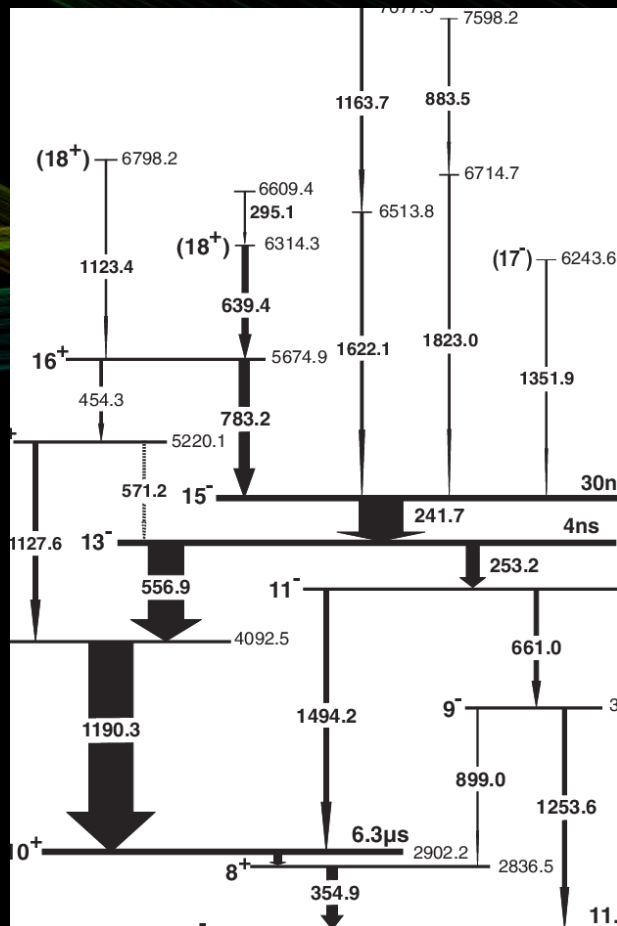
$$|I_1 - I_2| \equiv \Delta I \leq l \leq I_1 + I_2$$

Multipole	$E$	$M$
$l$	$\lambda (s^{-1})$	$\lambda (s^{-1})$
1	$1.03 \times 10^{14} A^{2/3} E_\gamma^3$	$3.15 \times 10^{13} E_\gamma^3$
2	$7.28 \times 10^7 A^{4/3} E_\gamma^5$	$2.24 \times 10^7 A^{4/3} E_\gamma^5$
3	$3.39 \times 10^1 A^2 E_\gamma^7$	$1.04 \times 10^1 A^{4/3} E_\gamma^7$
4	$1.07 \times 10^{-5} A^{8/3} E_\gamma^9$	$3.27 \times 10^{-6} A^2 E_\gamma^9$
5	$2.40 \times 10^{-12} A^{10/3} E_\gamma^{11}$	$7.36 \times 10^{-13} A^{8/3} E_\gamma^{11}$





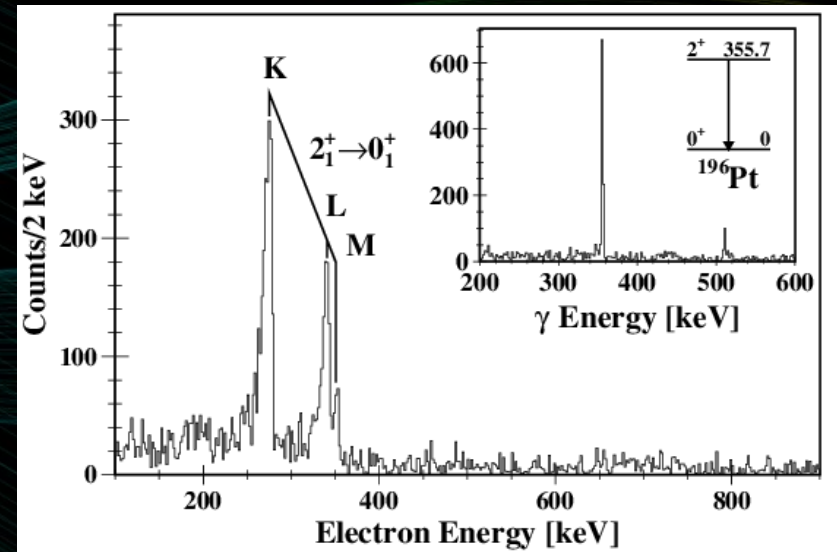
# gamma decay: level scheme



- arrows: thickness give the probability
- spin and parity, transition energies, level energy, branching ratios

# gamma decay

- Internal conversion: if gamma decay is forbidden (like isomeric states)
  - direct de-excitation by emitting a K electron
  - e.g.  $0^+ \rightarrow 0^+$  transition
  - internal conversion coefficients for K, L, M electrons: extracting multipolarity of the transitions!
  - $E_{e^-} = E_\gamma - E_{K(LM)}$



- Internal pair creation:
  - electron-positron pairs are created
  - $E_{e^-} + E_{e^+} = E_\gamma - 1022 \text{ keV}$

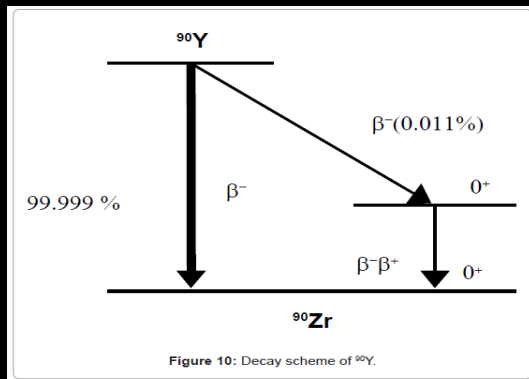
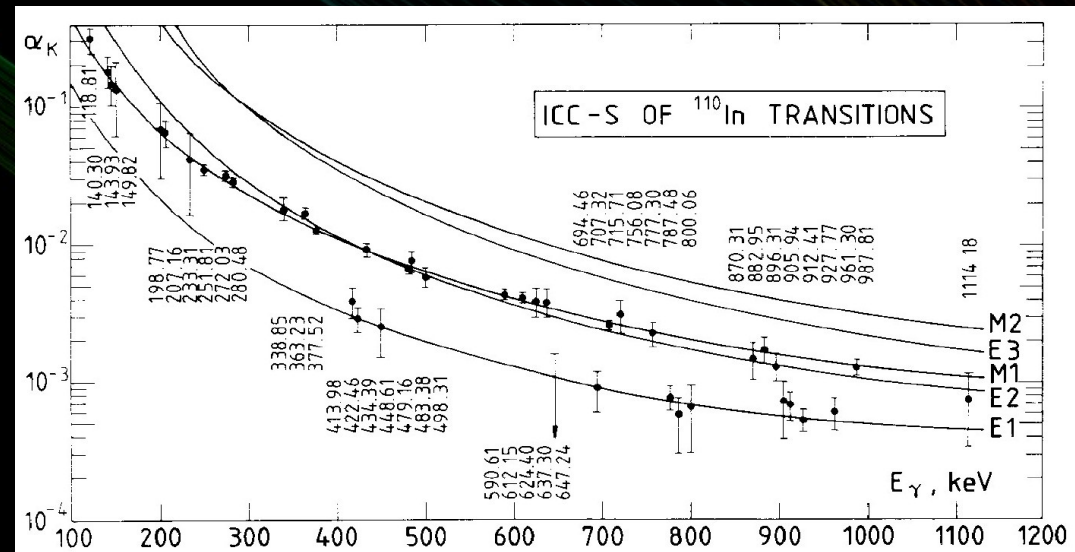
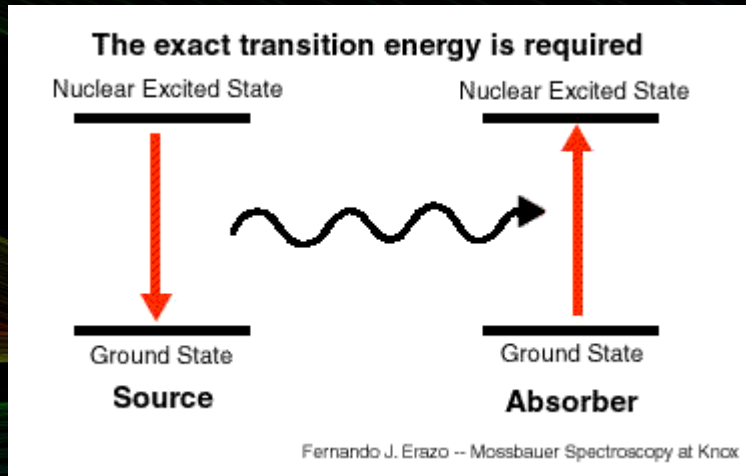


Figure 10: Decay scheme of  $^{90}\text{Y}$ .





# Mössbauer effect

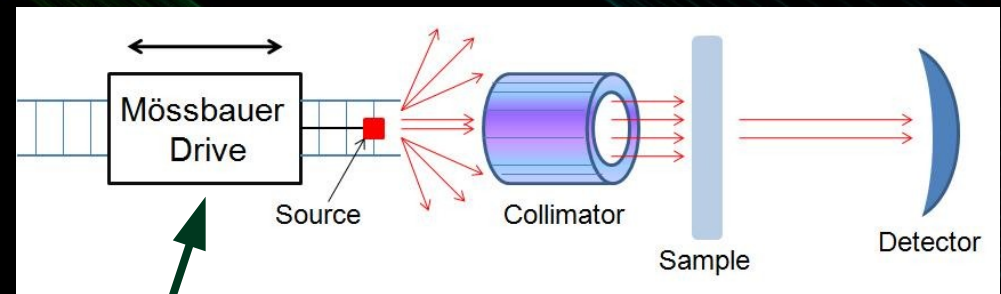


$$E_N = \frac{E_\gamma^2}{2 M_N c^2}$$

due to momentum conservation

- emitted  $E_\gamma < E_{\text{transition}}$  while excitation needs a bit larger  $E_\gamma \rightarrow$  if identical nuclei are in rest and free, no absorption can be observed due to recoil

- If nuclei are in a solid crystal  $\rightarrow$  phonons -vibrations of crystal lattice with discrete energy- can be emitted  $\rightarrow$  if no phonon, then absorption and emission is recoilless due to large ( $M$  momentum conservation with a crystal as a whole)
- „recoilless nuclear resonance fluorescence“
- Mössbauer spectroscopy:
  - source and sample material is the same (as crystal)
  - source is moving with velocity  $v \rightarrow$  Doppler shift



$$E_\gamma = E_{\gamma 0} \left( 1 + \frac{v}{c} \right)$$

# Mössbauer spectroscopy

- A typical experiment: changing (scanning) the velocity of source with a linear motor → and see the resonance absorption lines → Intensity vs. velocity is measured (typically  $\pm 10$ - $20$  mm/s)
- Energy resolution is extremely good: up to  $dE/E = 10^{-13}$  !!
- Isomer shift, quadrupole splitting and hyperfine splitting can be measured

